

## Chapter 11.

### Optical Processes and Excitons

$\epsilon(\omega, k)$  : electromagnetic field에 대한 반응  
 이것은 electronic band structure에 매우 sensitive 하다.

$$\epsilon = \epsilon' + i\epsilon''$$

Directly accessible function :  $R(\omega)$ ,  $n(\omega)$ ,  $k(\omega)$   
 extinction coefficient  
 refractive index  
 reflectance

Optical reflectance

$$\frac{E(\text{refl})}{E(\text{inc})} \equiv r(\omega) \equiv \rho(\omega) \exp[i\theta(\omega)] \quad \text{reflective coefficient라 불리운다.}$$

$$r(\omega) = \frac{n + ik - 1}{n + ik + 1}$$

definition  $n(\omega)$  and  $k(\omega)$  are related

$$\begin{aligned} \sqrt{\epsilon(\omega)} &\equiv n(\omega) + ik(\omega) \\ &\equiv N(\omega) \end{aligned}$$

$$E_y(\text{inc}) = E_{y_0} \exp[i(kx - \omega t)]$$

$$\begin{aligned} E_y(\text{trans}) &\propto \exp[i(n + ik)kx - \omega t] \\ &= \exp[-kx] \exp[i(nkx - \omega t)] \end{aligned}$$

Reflectance R

$$\begin{aligned} R &= E^*(\text{refl})E(\text{refl}) / E^*(\text{inc})E(\text{inc}) \\ &= r^* r \end{aligned}$$

$$\begin{aligned} &= \rho^2 \quad \rho = |r| \\ \epsilon(\omega) &= \epsilon'(\omega) + i\epsilon''(\omega) \end{aligned}$$

$$\begin{aligned} &= (n + ik)^2 \\ &= n^2 - k^2 + 2nki \end{aligned}$$

Kramers -Kronig Relations

Real을 알면 imaginary를 알 수 있고 imaginary를 알면 real을 알 수 있다.

$$\alpha(\omega) = \alpha'(\omega) + i\alpha''(\omega) \quad \text{Responsefunction}$$

$$X_\omega = \alpha(\omega)F_\omega$$

equation of motion

$$\sum_j M_j \left( \frac{d^2}{dt^2} + p_j \frac{d}{dt} + \omega_j^2 \right) X = F$$

$$\therefore \sum_j M_j (-\omega^2 - i\omega p_j + \omega_j^2) X_\omega = F_\omega$$

$$X = \int X_\omega \exp[-i\omega t] dt \quad F = \int F_\omega \exp[-i\omega t] dt$$

$$\begin{aligned} \therefore \alpha(\omega) &= \sum_j \frac{1}{M_j (\omega^2 - i\omega p_j + \omega_j^2)} \\ &= \sum_j \frac{f_j (-\omega^2 + i\omega p_j + \omega_j^2)}{(-\omega^2 - i\omega p_j + \omega_j^2)(-\omega^2 + i\omega p_j + \omega_j^2)} \\ &= \sum_j f_j \frac{(\omega_j^2 - \omega^2 + i\omega p_j)}{(\omega_j^2 - \omega^2)^2 + \omega^2 p_j^2} \end{aligned}$$

$f_j$

all positive for a passive system

$p_j$  relation frequency

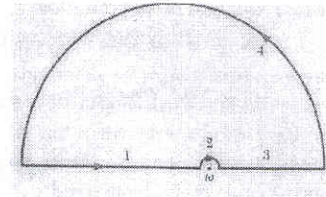
Kramers - Kronig relation

(1) The poles of  $\alpha(\omega)$  are all below the real axis

(2)  $\frac{\alpha(\omega)}{\omega} \xrightarrow{\omega \rightarrow \infty} \text{vanish}$

$\alpha'(\omega)$  is even

(3)  $\alpha''$  is odd



$$\alpha(\omega) = \frac{1}{\pi i} p \int_{-\infty}^{\infty} \frac{\alpha(s)}{s - \omega} ds$$

$$\alpha'(\omega) = \frac{1}{\pi} p \int_{-\infty}^{\infty} \frac{\alpha''(s)}{s - \omega} ds$$

$$= \frac{1}{\pi} p \left[ \int_0^{\infty} \frac{\alpha''(s)}{s - \omega} ds + \int_{-\infty}^0 \frac{\alpha''(p)}{p - \omega} dp \right]$$

$$= \frac{1}{\pi} p \left[ \int_0^{\infty} \frac{\alpha''(s)}{s - \omega} ds + \int_{\infty}^0 \frac{\alpha''(-s)}{-s - \omega} d(-s) \right]$$

$$= \frac{1}{\pi} p \left[ \int_0^{\infty} \frac{\alpha''(s)}{s - \omega} ds + \int_0^{\infty} \frac{\alpha''(s)}{s + \omega} ds \right]$$

$$= \frac{1}{\pi} p \int_0^{\infty} \frac{2s\alpha''(s)}{s^2 - \omega^2} ds$$

$$\alpha'(\omega) = \frac{2}{\pi} p \int_0^{\infty} \frac{s\alpha''(s)}{s^2 - \omega^2} ds$$

$$\int_{\omega}^{\infty} \frac{\alpha''(s)}{s + \omega} ds$$

$$\alpha''(\omega) = -\frac{1}{\pi} p \int_{-\infty}^{\infty} \frac{\alpha'(s) ds}{s-\omega}$$

$$= -\frac{2\omega}{\pi} p \int_0^{\infty} \frac{\alpha'(s) ds}{s^2-\omega^2}$$

Kramers - Kronig Relation 을 적용하자.

$$\ln r(\omega) = \ln R^{\frac{1}{2}}(\omega) + i\theta(\omega)$$

$$\theta(\omega) = -\frac{\omega}{\pi} p \int_0^{\infty} \frac{\ln R(s)}{s^2-\omega^2} ds$$

$$= -\frac{i}{2\pi} p \int_0^{\infty} \ln \left| \frac{s-\omega}{s+\omega} \right| \frac{d \ln R(s)}{ds} ds$$

$$\frac{1}{s^2-\omega^2} = \frac{1}{(s-\omega)(s+\omega)}$$

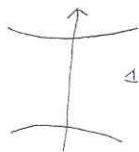
$$= \left( \frac{1}{s-\omega} - \frac{1}{s+\omega} \right) \frac{1}{2\omega}$$

Mathematical note

$$\alpha(\omega) = \sum_j \frac{(\omega_j^2 - \omega^2) + i\omega p_j}{(\omega_j^2 - \omega^2) + \omega^2 P_j^2}$$

### Electronic

Interband transition



1. 이렇게 gap보다 더 크면 optical spectroscopy는 broad feature만 보인다.

direct interband absorption은  $\hbar\omega = \epsilon_c(\vec{k}) - \epsilon_v(\vec{k})$ .

### Excitons

bound electron - hole pair를 exciton이라 한다.

exciton이 crystal에서 움직이면서 에너지를 전달하다.

전기는 전달 안 함

Exciton은 모든 insulating crystal에서 forming 될 수 있다.

gap이 indirect이면 unstable 하다.

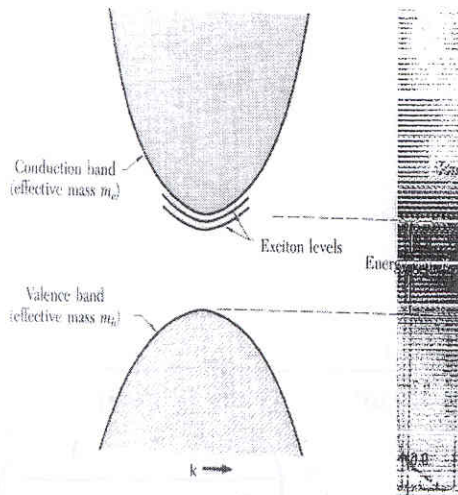
모든 exciton이 다 unstable하다. - 전자가 hole로 들어간다.

threshold of this process

$\hbar\omega > E_g$  in a direct process

$\hbar\omega > E_g - \hbar\Omega$  indirect process

정확하게는 exciton의 binding 에너지를 고려해야 한다.

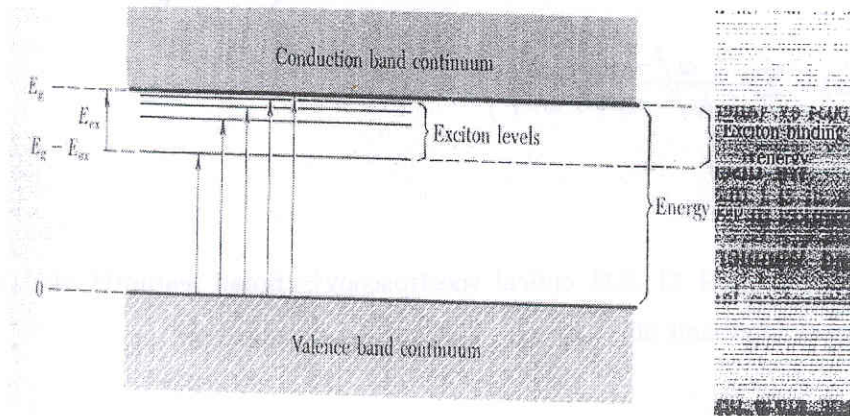


binding energy :

1. exciton을 만드는 에너지와 free electron, free hole을 만드는 에너지의 차이
2. recombination luminescence , free electron - hole combination
3. photo - ionization of excitons

⇒ exciton is weakly bounded: Mott and Wannier

⇒ tightly bounded : Frenkel



Frenkel Excitons

Tightly bound.

⇒ hole은 바로 옆에 있다.( single atom에 있다.)

⇒ 이 exciton이 hopping을 한다.

⇒ Crystalline inert gas : Frenkel model로 설명

Atomic crypton : lowest strong atomic transition at 9.99eV

Crystal crypton :

10.2 eV (gap) 이 11.7

유전자  
Crypton 9.99eV  
결정내에서는 10.2eV

exciton의 ground state energy

$$11.7 - 10.2 = 1.5 \text{ eV (exciton의 ground state energy)}$$

Exciton의 propagation

옆으로 옆으로 wave propagation

$$\Psi_g = u_1 u_2 \dots u_{N-1} u_N$$

single atom이 excited 되어 있다.

$$\varphi_j = u_1 u_2 \dots u_{j-1} u_{j+1} \dots u_N$$

$$H \varphi_j = \varepsilon \varphi_j + T(\varphi_{j-1} + \varphi_{j+1})$$

$\varepsilon$  : free atom excitation energy

$T$  : excitation  $j \rightarrow j-1$

$j \rightarrow j+1$

이 solution은 wave 모양이다.



$$\Psi_k = \sum \exp[ijka] \varphi_j$$

$$H \Psi_k = \sum_j \exp[ijka] H \varphi_j$$

$$= \sum_j \exp[ijka] [\varepsilon \varphi_j + T(\varphi_{j-1} + \varphi_{j+1})]$$

rearrange 하자

$$H \Psi_k = \sum \exp[ijka] [\varepsilon + T(\exp[ika] + \exp[-ika])] \varphi_j$$

$$= [\varepsilon + 2T \cdot \cos ka] \Psi_k$$

$$E_k = \varepsilon + 2T \cos ka$$

periodic boundary condition

$$K = \frac{2\pi s}{Na} \quad s = -\frac{1}{2}N, -\frac{N}{2} + 1, \dots, \frac{1}{2}N - 1$$

Alkali Halide : localized exciton on the negative halogenion

visible 에서는 exciton 안 보임

transparent

UV에서 exciton 보이고

Molecular crystals :

Van der Waals binding between molecules

covalent 가 매우 강하다.

이러한 exciton은 Frenkel excitons

Weakly bound ( Mott - Wannier )

$$U(r) = -\frac{e^2}{\varepsilon r}$$

에너지가 cond. 보다 아래에서 bound state 있다.

$$E_n = E_g - \frac{\mu e^4}{2\hbar^2 \varepsilon^2 n^2}$$

$$\frac{1}{\mu} = \frac{1}{m_e} + \frac{1}{m_h}$$

예  $Cu_2O \quad v(cm^{-1}) = 17508 - \frac{800}{n^2}$

Exciton Condensation into Electron - Hole drop

electron - hole drop 이 Ge 속에서 생긴다.

초전도연구단  
단장 이성익

$\hbar \omega > E_g$  이면 free - electron and free hole이 생긴다.

recombination  $8\mu s$

electron - hole recombination and annihilation

$8\mu s$

만약 exciton concentration이 high이면 drop의 life time은  $40\mu s$

strained germanium  $600\mu s$

excitons의 drop은 degenerated Fermi gas가 되고 metallic properties가 나타난다.

### Raman effect

Photon inelastically scatter

creation or annihilation of a phonon or magnon

Selection rule

$$\omega = \omega' \pm \Omega ; \quad k = k' \pm K$$

$\Omega, K$  : phonon의 creation 또는 annihilation

2nd order Raman effect :

two phonons are involved in the inelastic scattering of photon

strain dependence of electronic polarizability

$$\alpha = \alpha_0 + \alpha_1 u + \alpha_2 u^2 + \dots$$

$\alpha_1$  : phonon amplitude

$$u(t) = u_0 \cos \Omega t$$

induced electric dipole moment

$$\alpha_1 E \cdot u$$

$$= \alpha_1 E_0 u_0 \cos \omega t \cos \Omega t$$

$$= \frac{1}{2} \alpha_1 E_0 u_0 [\cos(\omega + \Omega)t + \cos(\omega - \Omega)t]$$

$\omega - \Omega$  : stokes line

$\omega + \Omega$  : anti stokes line

$$I(\omega - \Omega) \propto | \langle n_{K+1} | u | n_K \rangle |^2 \propto n_{K+1}$$

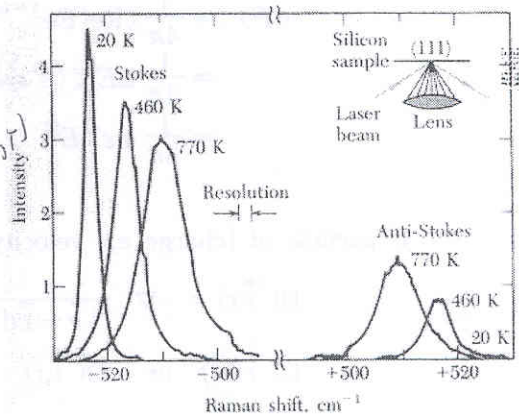
$$I(\omega + \Omega) \propto | \langle n_{K-1} | u | n_K \rangle |^2 \propto n_K$$

Population

$$\frac{I(\omega + \Omega)}{I(\omega - \Omega)} = \frac{\langle n_K \rangle}{\langle n_K \rangle + 1} = \exp[-\hbar \Omega / k_B T]$$

$$\therefore \langle n_K \rangle = \frac{1}{\exp[\hbar \omega / k_B T] - 1}$$

$$\langle n_{K+1} \rangle = \frac{1 + \exp(\hbar \omega / k_B T) - 1}{\exp(\hbar \omega / k_B T) - 1} = \frac{\exp(\hbar \omega / k_B T)}{\exp(\hbar \omega / k_B T) - 1}$$



Electron Spectroscopy with X-ray

X-ray Photoemission from solids (XPS)

Ultraviolet Photoemission (UPS)

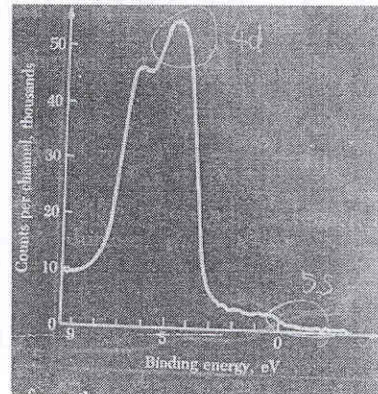
Density of state(  $D(\epsilon)$  )

photon이 incident photon - binding energy = K.E.

전자는 50Å 정도 나올 수 있다.

ESCA : resolution 0.6 eV

Electron Spectroscopy for Chemical  
valence structure of silver =>그림  
Analysis



Excitation

99.2 eV binding energy

117 eV에서 또 보인다.

single plasmon excitation

134.7 eV : double plasmon excitation

Energy loss of fast particles in a solid

전자를 빛 대신쓰자.

$Im\{\epsilon(\omega)\}$  : energy loss by an electromagnetic wave in a solid

$Im\left\{\frac{1}{\epsilon(\omega)}\right\}$  : energy loss by a charged particle that penetrated a solid

$$P = \frac{1}{4\pi} E \cdot \frac{\partial D}{\partial t}$$

$$= \left[ \frac{1}{4\pi} E \cdot (-i\omega) \right] \cdot [\epsilon(\omega) E e^{-i\omega t}]$$

$$E = E_0 e^{-i\omega t}$$

$$D = D_0 e^{-i\omega t}$$

$$\begin{aligned}
 \langle P \rangle &= \frac{1}{4\pi} \langle \text{Re}\{E e^{-i\omega t}\} \cdot \text{Re}\{-i\omega D e^{-i\omega t}\}\rangle \\
 &= \frac{1}{4\pi} \omega E^2 \langle (\epsilon' \sin \omega t + \epsilon'' \cos \omega t) \cos \omega t \rangle \\
 &= \frac{1}{8\pi} \omega \epsilon'' E^2
 \end{aligned}$$

$$\begin{aligned}
 &-i\omega (\epsilon' + i\epsilon'') (\cos \omega t + i \sin \omega t) \\
 &= \left[ \epsilon' \cos \omega t - \epsilon'' \sin \omega t \right] \\
 &\quad + i (\epsilon' \sin \omega t + \epsilon'' \cos \omega t)
 \end{aligned}$$

If particle of [charge  $e$ , velocity  $v$ ] enters a crystal:

$$D(\vec{r}, t) = -\nabla \frac{e}{|\vec{r} - v\vec{t}|}$$

$$D(\vec{r}, \omega) \quad \text{or} \quad D(\vec{r}, t)$$

$$E(\omega, \vec{k}) = D(\omega, \vec{k}) / \epsilon(\omega, \vec{k})$$

Time average power dissipation

$$\begin{aligned}
 P(\omega, \vec{k}) &= \frac{1}{4\pi} \langle \text{Re} \left\{ \frac{D(\omega, \vec{k})}{\epsilon(\omega, \vec{k})} e^{-i\omega t} \right\} \cdot \text{Re}\{-i\omega D(\omega, \vec{k}) e^{-i\omega t}\} \rangle \\
 &= \frac{1}{4\pi} \omega D^2(\omega, \vec{k}) \langle \left[ \left( \frac{1}{\epsilon} \right)' \cos \omega t + \left( \frac{1}{\epsilon} \right)'' \sin \omega t \right] [-\sin \omega t] \rangle
 \end{aligned}$$

Whence

$$\begin{aligned}
 P(\omega, \vec{k}) &= \frac{-1}{8\pi} \omega \left( \frac{1}{\epsilon} \right)'' D^2(\omega, \vec{k}) \\
 &= \frac{1}{8\pi} \omega \frac{\epsilon''(\omega, \vec{k})}{|\epsilon|^2} D^2(\omega, \vec{k})
 \end{aligned}$$

$$\text{energy loss function} \quad \propto \text{Im} \left( \frac{1}{\epsilon(\omega, \vec{k})} \right)$$

위  $P(\omega, k)$ 가  $k$ 에 무관하면

$$P(\omega) = -\frac{2}{\pi} \cdot \frac{e^2}{\hbar v} \text{Im} \left\{ \frac{1}{\epsilon} \right\} \ln \left( \frac{k_0 v}{\omega} \right)$$