

Manipulating Spins in Topological Insulators

Joon Sue Lee

California NanoSystems Institute, UC Santa Barbara

Korea University, Sep 8, 2016



C-SPIN



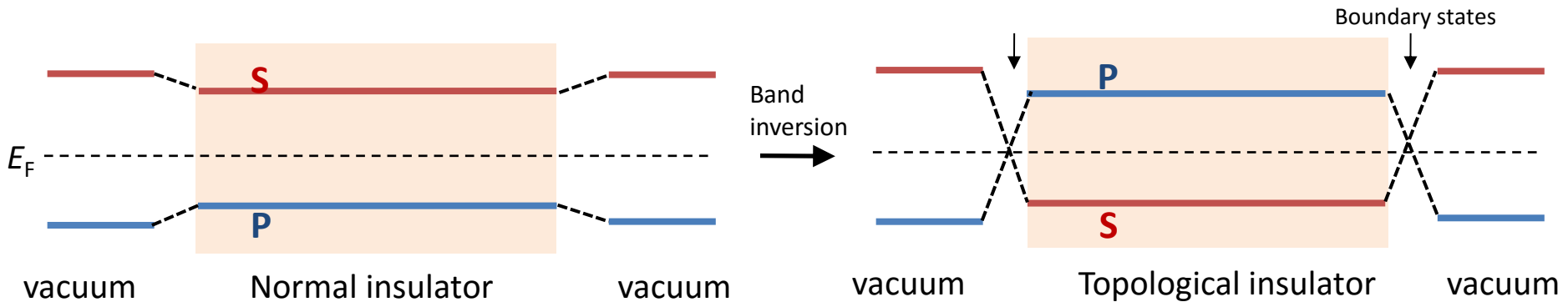
Introduction to topological insulators

Topology: mathematical study of the properties that remain under continuous deformations.



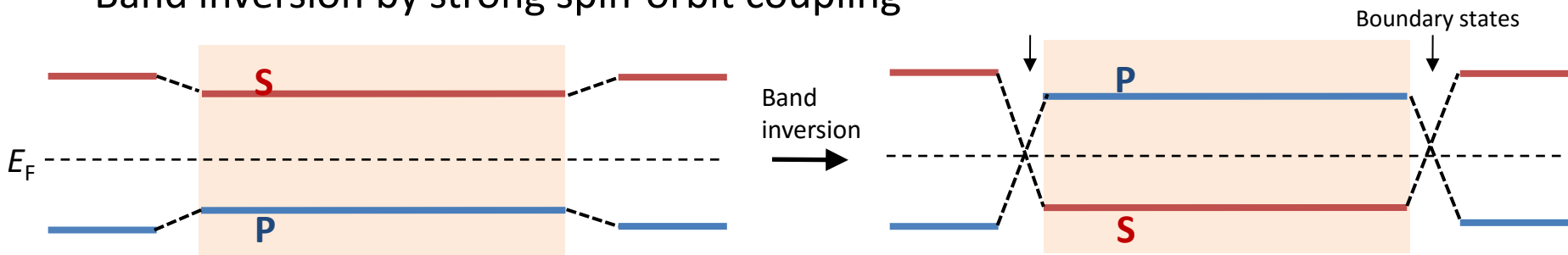
Graphic: www.pinterest.com

In condensed matter physics, continuous deformation \rightarrow change without closing a bulk energy gap.



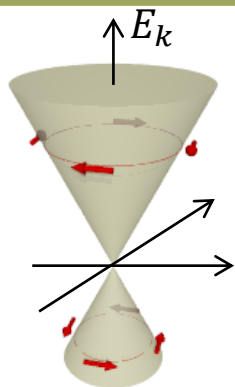
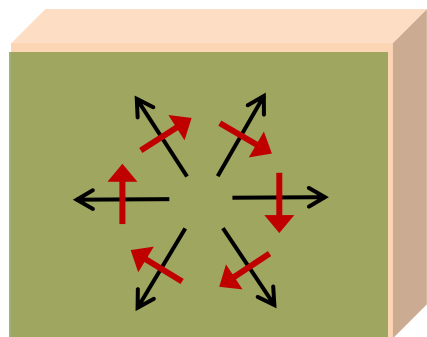
Introduction to topological insulators

- Band inversion by strong spin-orbit coupling



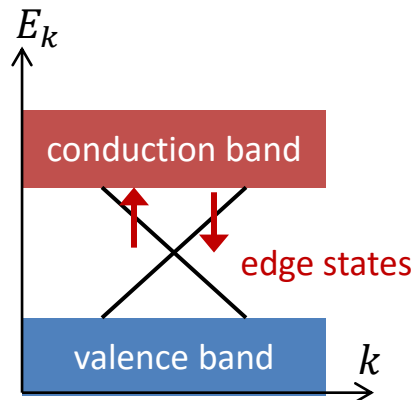
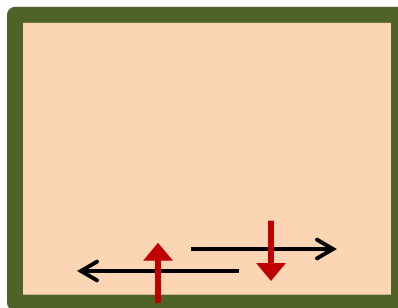
3D topological insulator

Surface states



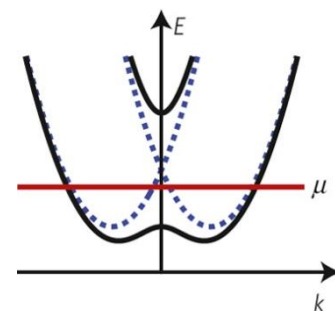
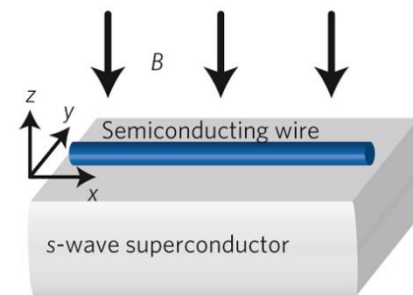
2D topological insulator

Edge states



1D topological superconductor

End states
(Majorana zero modes)



Hasan & Kane, *Rev. Mod. Phys.* (2010)

Qi & Zhang, *Physics Today* (2010)

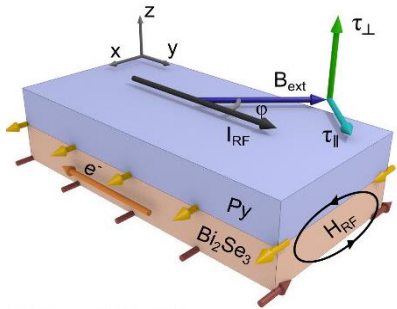
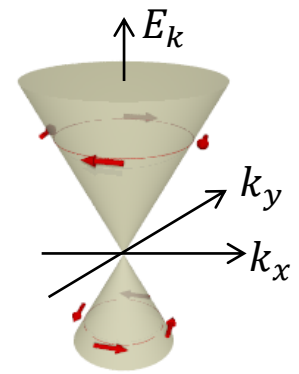
Moore, *Nature* (2010)

Qi & Zhang, *Rev. Mod. Phys.* (2011)

Moore & Balents, *Phys. Rev. B* (2007)

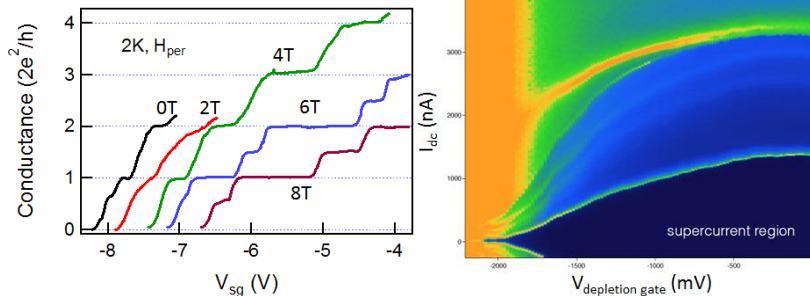
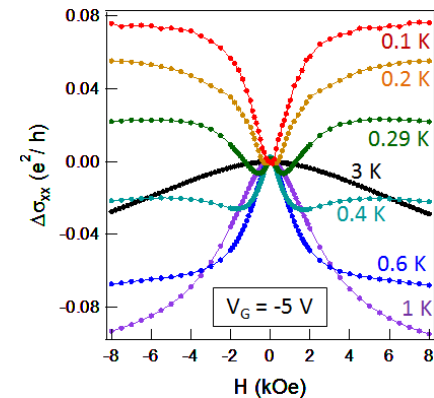
Outline

1. Spin-polarized surface states in 3D topological insulators
 - Electrical detection of spin polarization



2. Spin-charge conversion in topological insulators
 - Charge-to-spin conversion: spin transfer torque
 - Spin-to-charge conversion: spin pumping

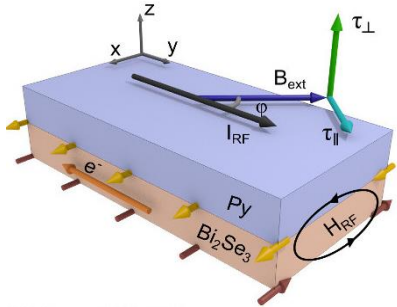
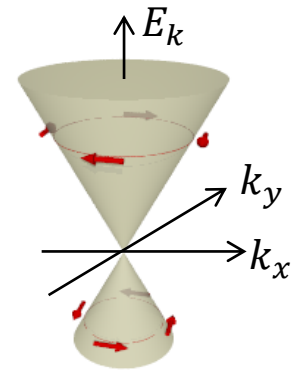
3. Modified surface states by breaking time-reversal symmetry
 - Interfacing a ferromagnetic insulator to a TI



4. Topological superconductivity towards quantum computation
 - 1D transport and superconducting proximity effect in epi-Al/2DEG

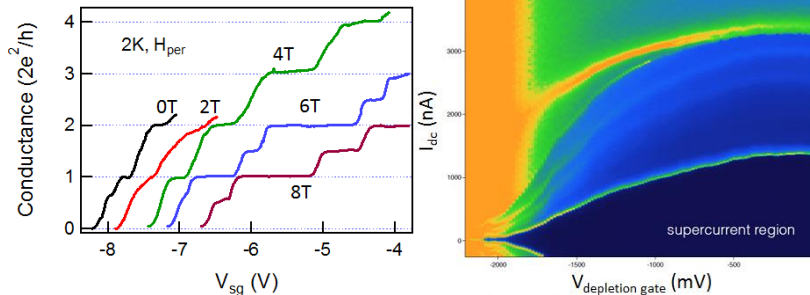
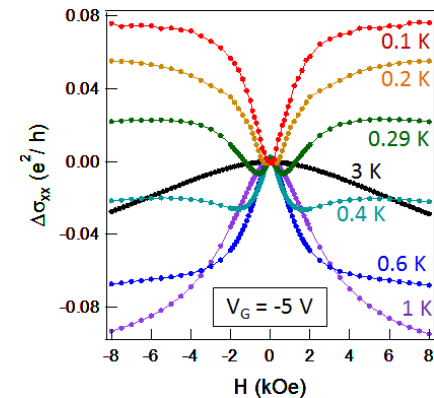
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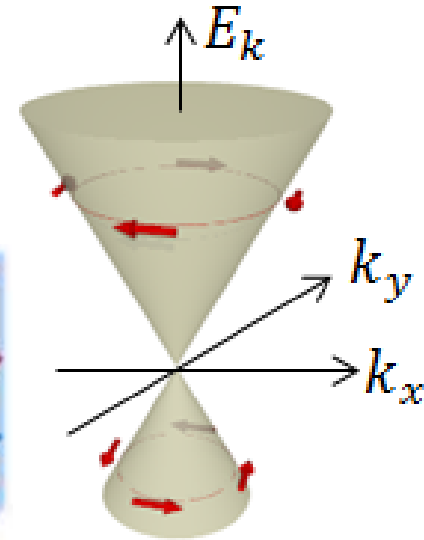
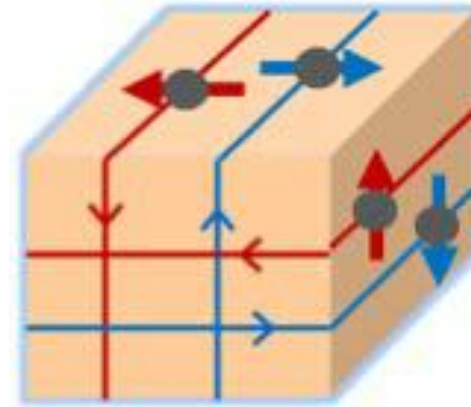
3D topological insulators

A state of matter that has an **insulating bulk** and **spin-dependent surface states** due to **spin-orbit coupling**.

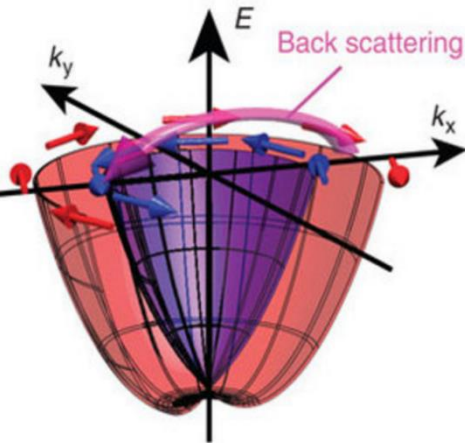
Surface states

$$H_{surf} = \hbar v_F (\sigma^x k_y - \sigma^y k_x)$$

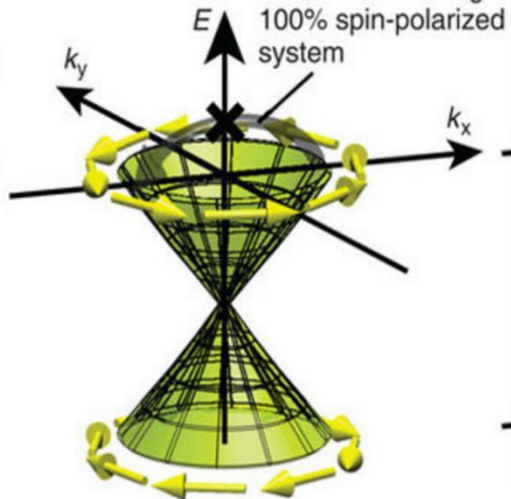
- Spin-momentum locking
- Protected by time-reversal symmetry
- Robust against back-scattering



Rashba states



Surface states



Time-reversal switches both the direction of electron motion and the direction of spin.

Time-reversal invariance:

$$TH(k)T^{-1} = H(-k)$$

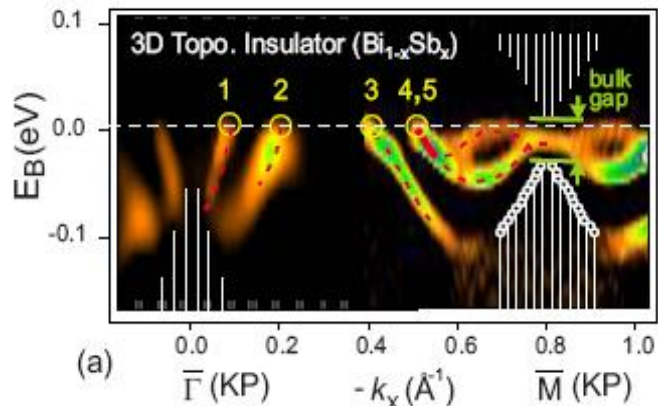
$$TH(k)T^{-1} = \hbar v_F [(-\sigma^x)k_y - (-\sigma^y)k_x]$$

$$H(-k) = \hbar v_F [\sigma^x(-k_y) - \sigma^y(-k_x)]$$

Invariant!

3D topological insulators

$\text{Bi}_{1-x}\text{Sb}_x$: the first 3D topological insulator

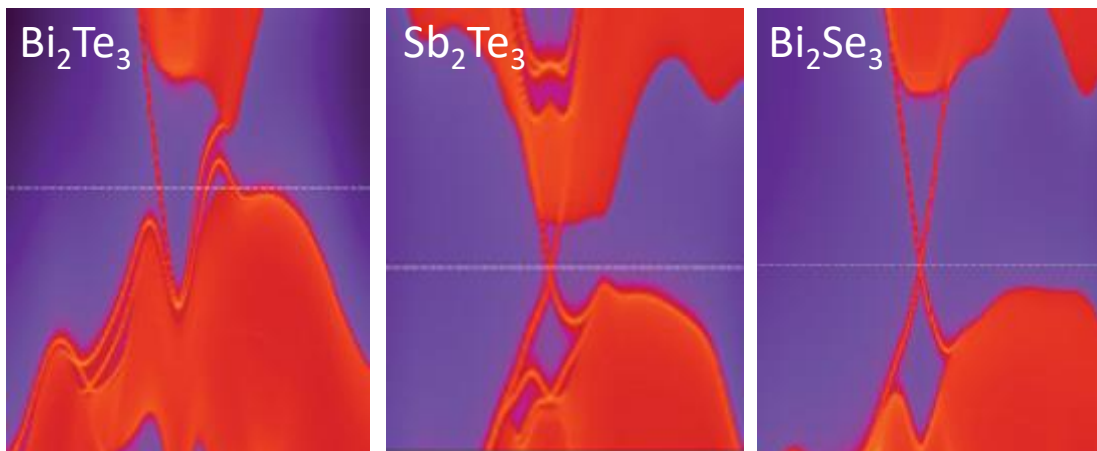


Prediction: Fu & Kane, *Phys. Rev. B* **76**, 045320 (2007).

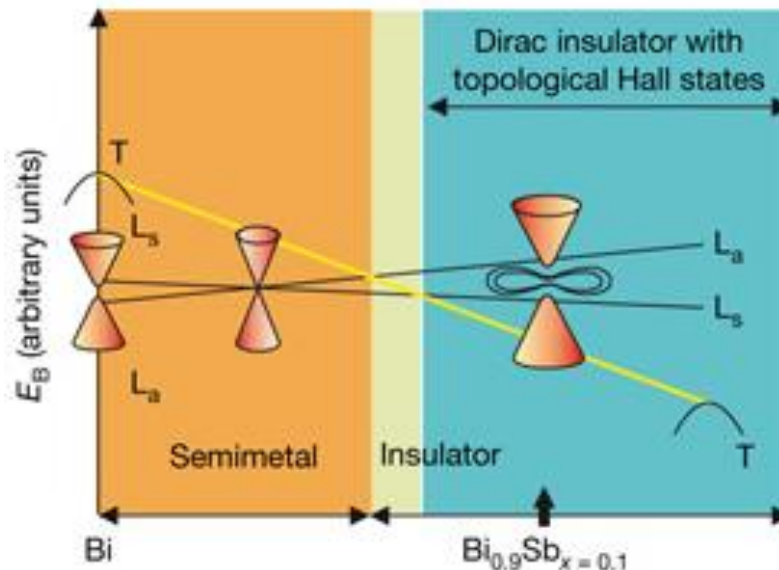
Experiment: Hsieh, *et al.*, *Nature* **452** (2008)

Bi-chalcogenides with a single Dirac cone

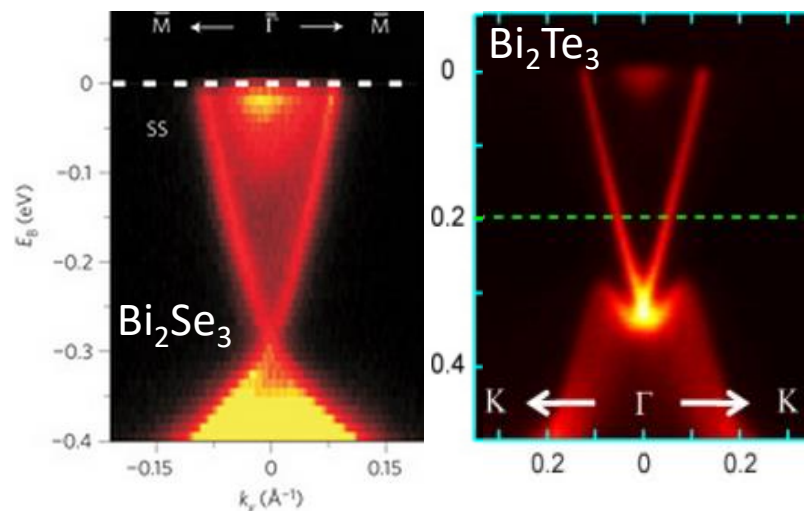
<First-principles electronic structure calculations>



Prediction: Zhang, *et al.*, *Nat. Phys.* **5**, 438 (2008)



<Angle-resolved photoemission spectroscopy>

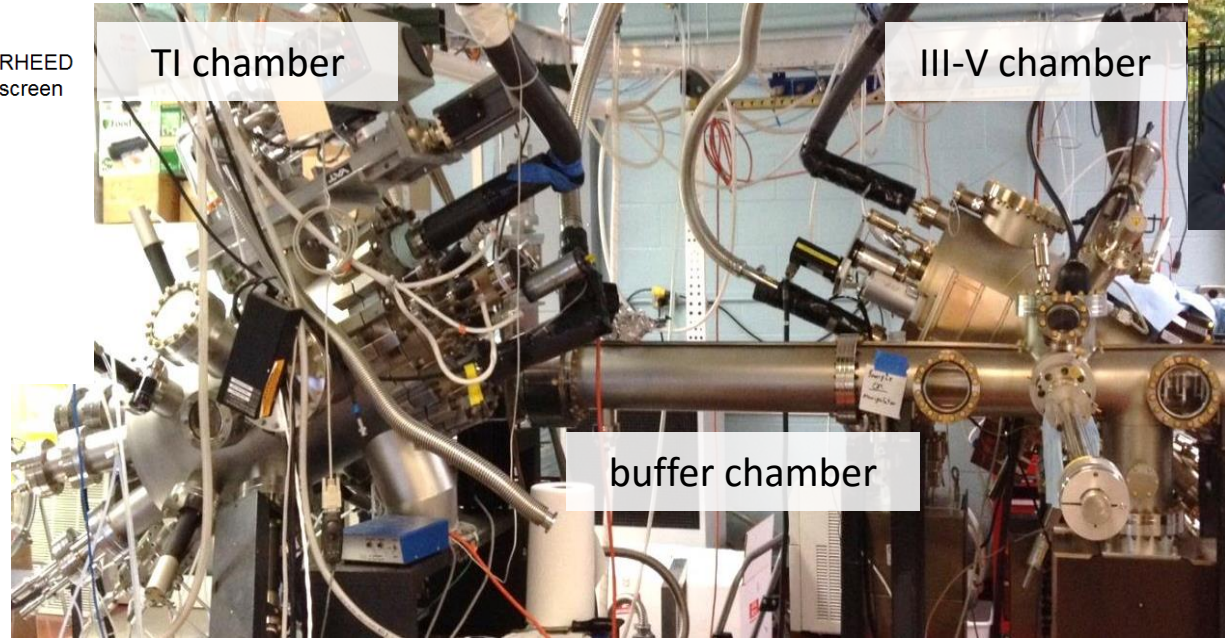
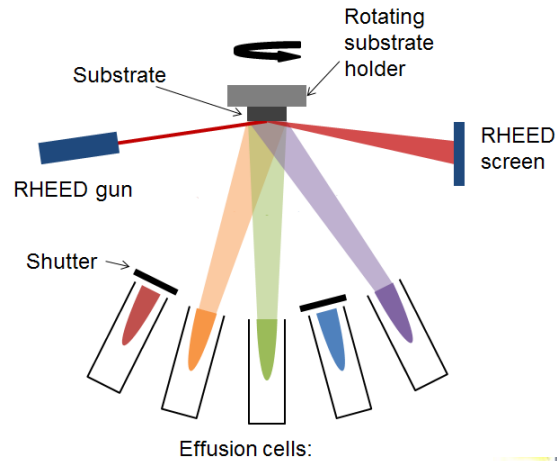


Experiment: Xia, *et al.*, *Nat. Phys.* **5**, 398 (2009)

Chen, *et al.*, *Science* **325**, 178 (2009) 7

Topological insulator thin films by MBE

MBE systems at Penn State (Nitin Samarth group)



Bi, Sb, Se, Te, Mn, and Cr.
(e-beam sources: Fe, Co, Nb, and V)

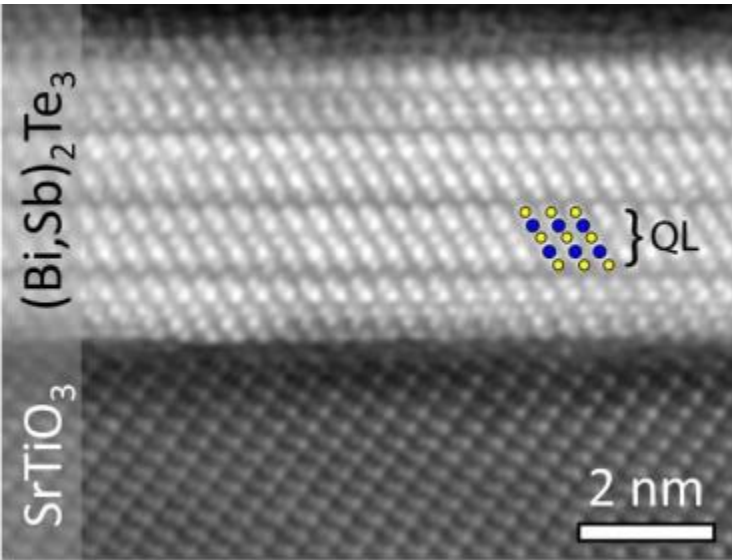
Ga, Al, In, As, Sb, Si, Be and Mn.

Molecular beam epitaxy (MBE) of topological insulators

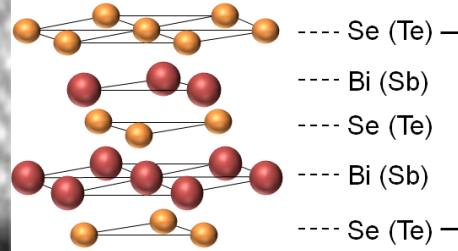
- Bi-chalcogenides: Bi_2Se_3 , Bi_2Te_3 , Sb_2Te_3 and $(\text{Bi,Sb})_2(\text{Te,Se})_3$...
- Magnetic doping: Cr, Mn, and V
- On various substrates: InP, GaAs, ZnSe, Si, sapphire, diamond, SrTiO_3 , graphene, h-BN, NbSe₂, YIG, LuIG, BaM, EuS, BSSCO, ...

Topological insulator thin films by MBE

Cross-sectional STEM

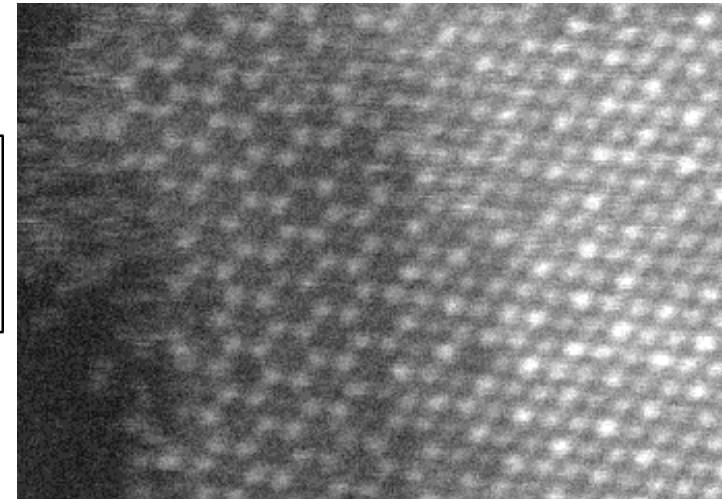


Rhombohedral crystal structure

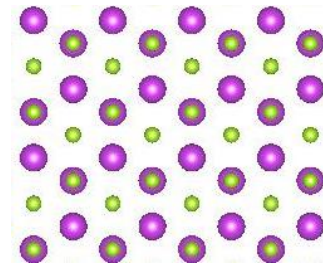


- Quintuple layer (QL)
- van der Waals bonding between QLs

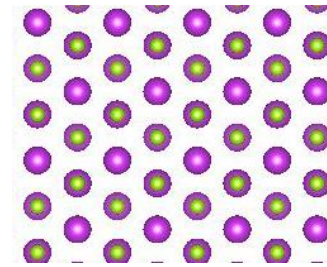
Plan-view STEM of $(\text{Bi,Sb})_2\text{Te}_3/\text{h-BN}$



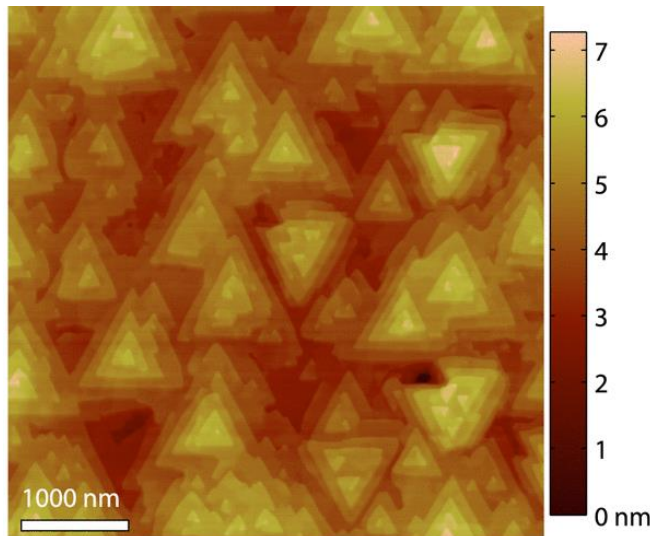
1 QL



multiple QLs



AFM of Bi_2Se_3 on InP

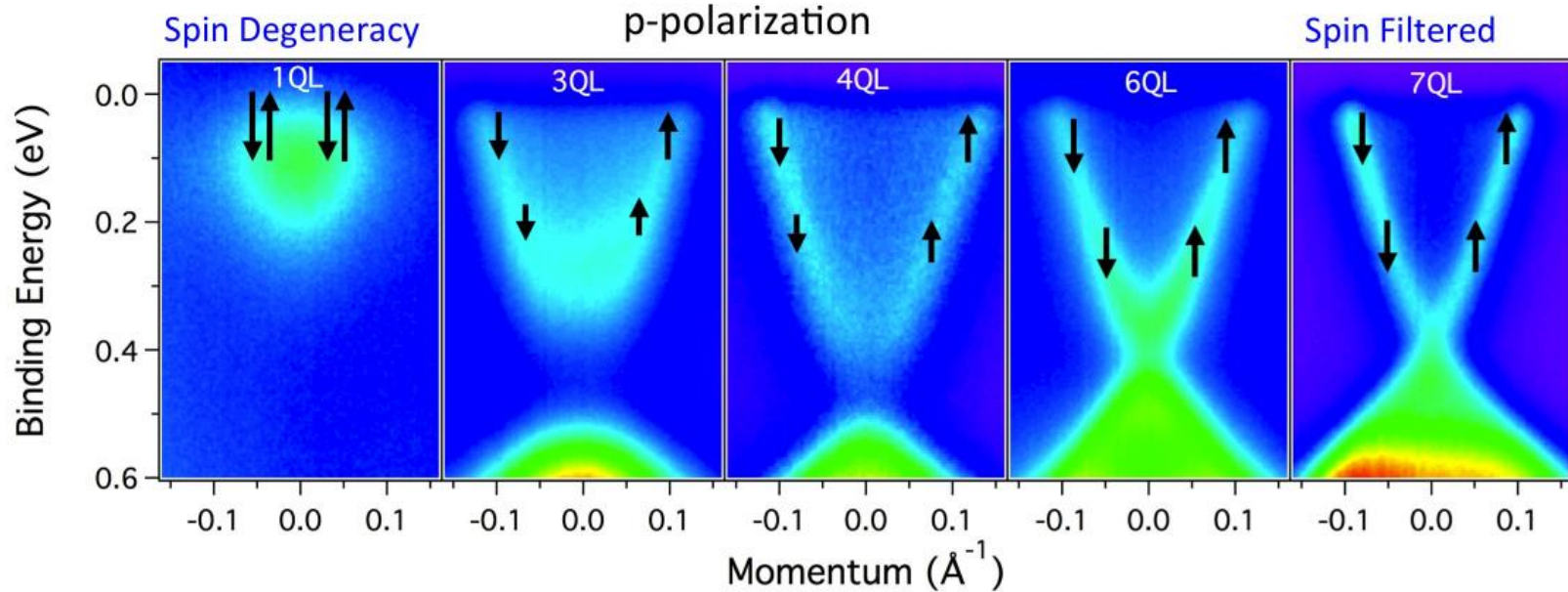


Film quality is limited by defects, such as vacancies, dislocations, twin boundaries, and large grains.

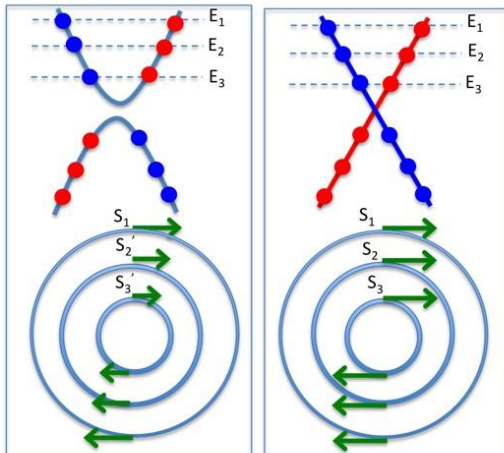
- Typical density: $1 \times 10^{17} - 1 \times 10^{19} \text{ cm}^{-3}$
- Typical electron mobility $< 2,000 \text{ cm}^2/\text{Vs}$

Spin texture in 3D topological insulators

Spin-resolved ARPES in 1-7 QL Bi_2Se_3 thin films show an evolution of the spin-polarized Dirac surface state.



M. Neupane *et al.*, *Nat. Comm.* **5**, 3841 (2014)

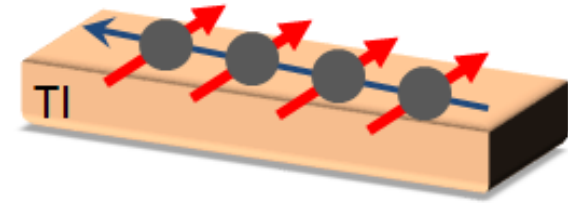


- A gap by hybridization between top and bottom surfaces (< 6 QL).
- The inherent spin-polarized surface state motivates for developing spin channel transport devices.

Topological spintronics

Spin polarized surface states in 3D topological insulators

→ Motivation for developing spintronic device applications



Spintronics (spin-based electronics)

: a field of electronics by control of electron spins for manipulation, storage, and transfer of information.

Spintronic effects:

- GMR → 2007 Nobel prize
- TMR
- Spin Hall effect
- Spin transfer torque
- Spin pumping
- Exchange bias with antiferromagnet



MRAM



Hard disc drive

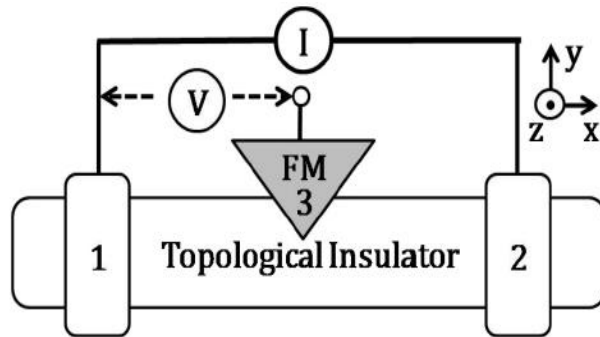
How to exploit the spin polarization in topological insulators?

Electrical detection of spin polarization in TIs

Experimental challenges:

- How to “electrically” measure the spin polarization in topological insulators?

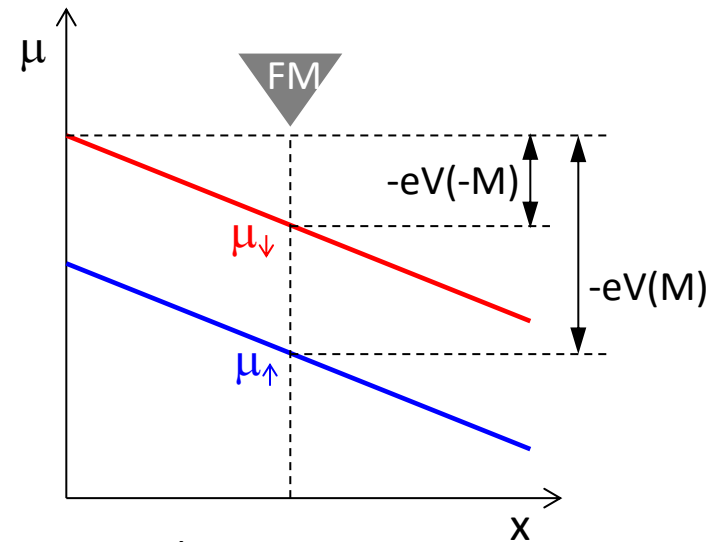
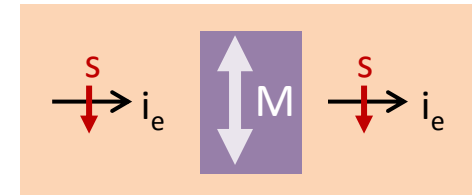
Theoretical proposal:
potentiometric geometry



Hong *et al.*, *Phys. Rev. B.* **86**, 085131 (2012)

$$\left[V(\vec{M}) - V(-\vec{M}) \right] / I = R_B(\vec{p} \cdot \vec{m})$$

ΔV is proportional to the projection of the TI spin polarization onto FM.



$$(\mu_{\uparrow} - \mu_{\downarrow})$$

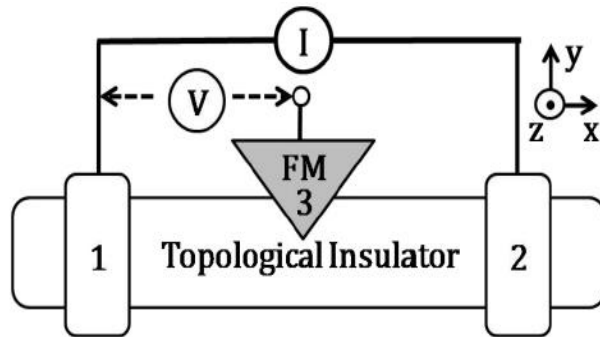
: spin electrochemical potential difference

Electrical detection of spin polarization in TIs

Experimental challenges:

- How to “electrically” measure the spin polarization in topological insulators?

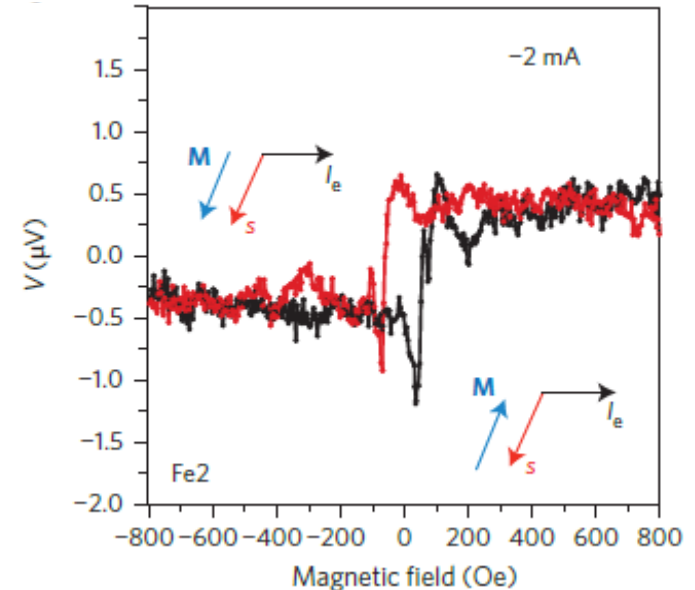
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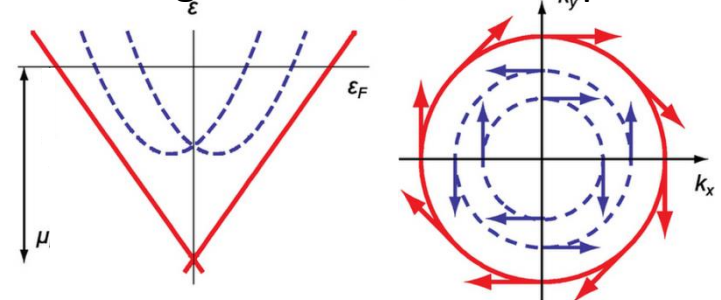
Li, *et al.*, *Nature Nano.* **9**, 218 (2014)

Tang *et al.*, *Nano Lett* **14**, 5423 (2014)

Ando *et al.*, *Nano Lett.* **14**, 6226 (2014)

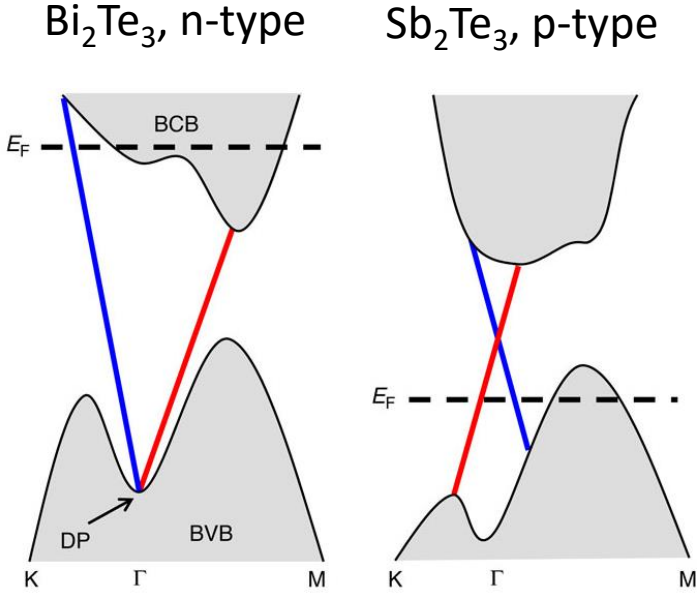
Tian *et al.*, *Sci. Reports* **5**, 14293 (2015)

→ Measuring at a fixed chemical potential.



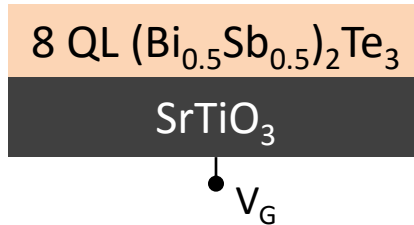
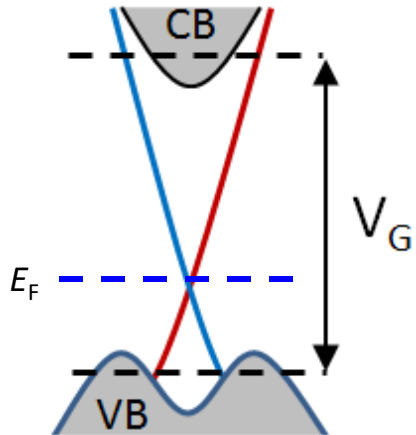
Can we measure as we tune the chemical potential?

Ambipolar transport in a topological insulator film

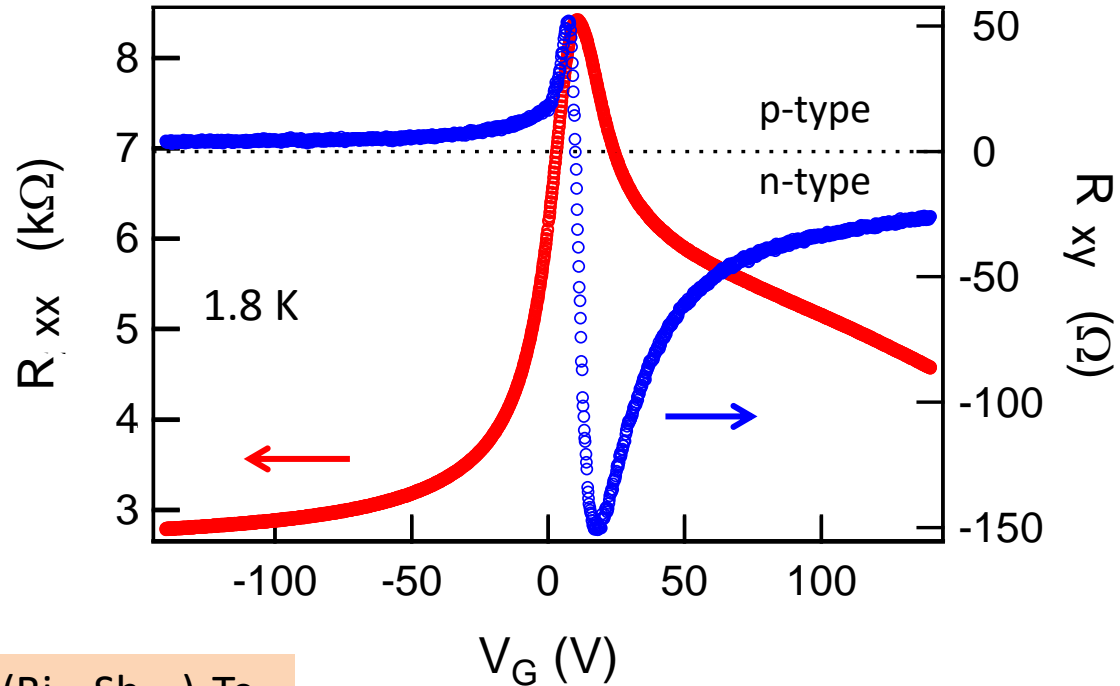
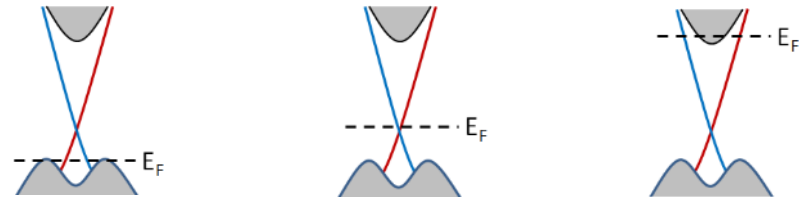


Zhang, *et al.*, *Nat. Commun.* 2, 574 (2011)

(Bi,Sb)₂Te₃
- E_F near the Dirac point

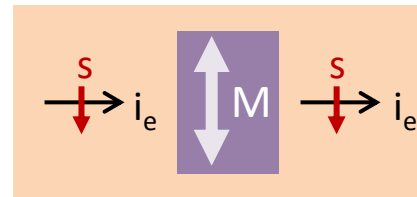
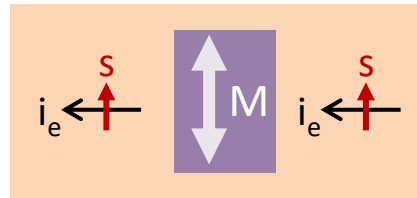
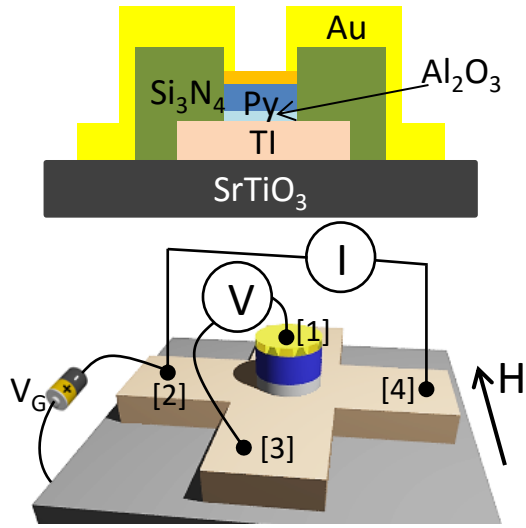


V_G tunes the TI chemical potential through Dirac point.

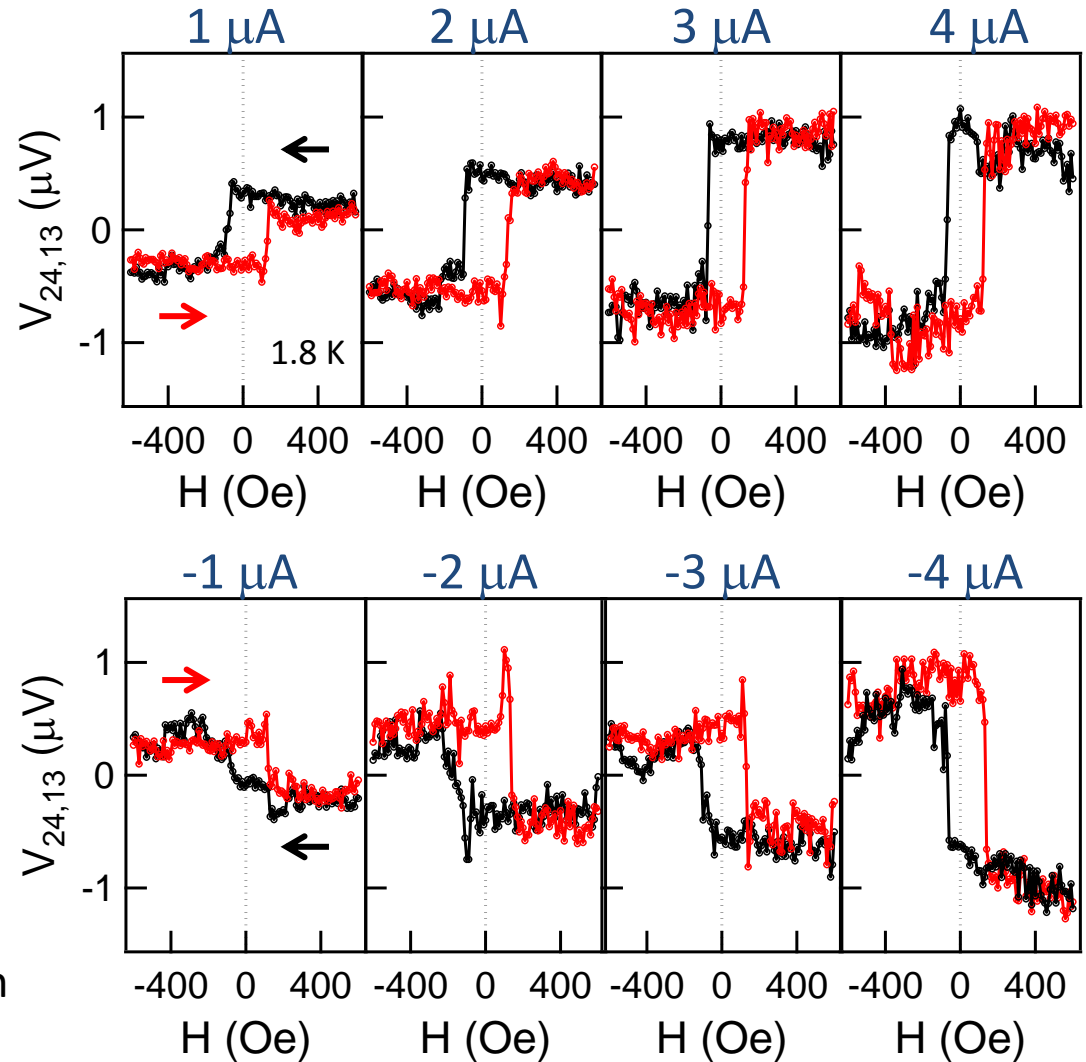


Electrical detection of spin polarization

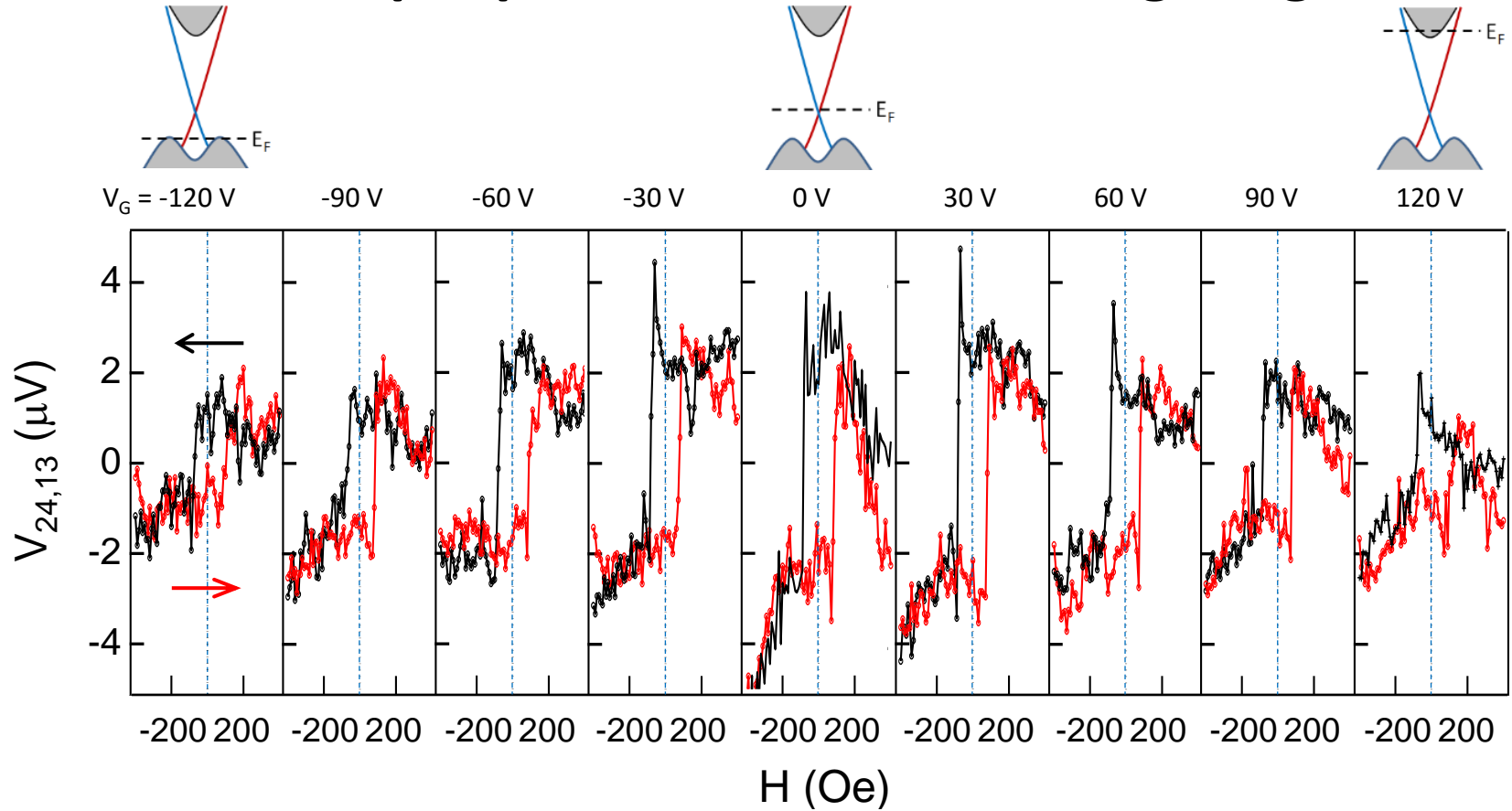
Lee et al., *Phys. Rev. B* **92**, 155312 (2015)



Hysteretic spin signal associated with TI spin polarization and Py switching.

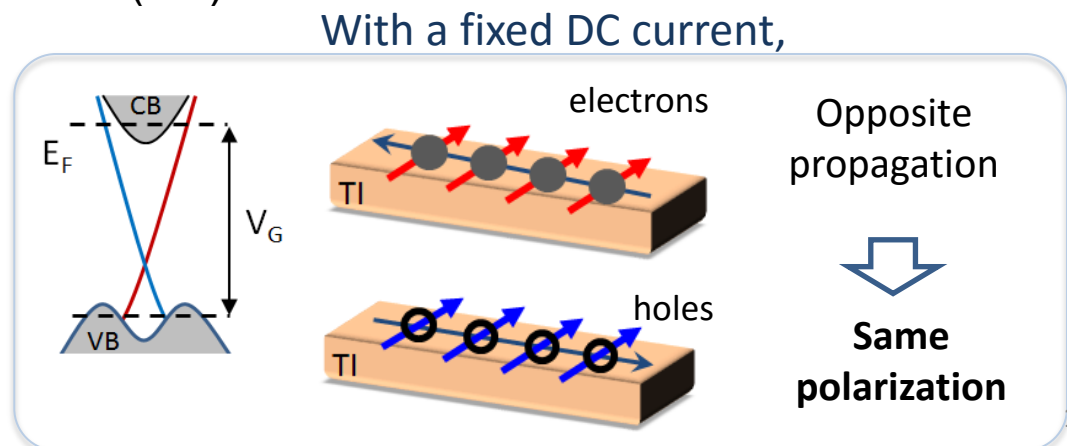


Detection of spin polarization with back gating



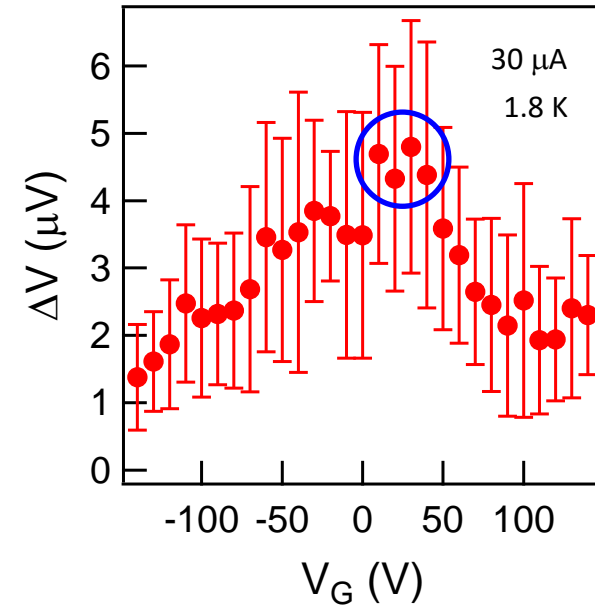
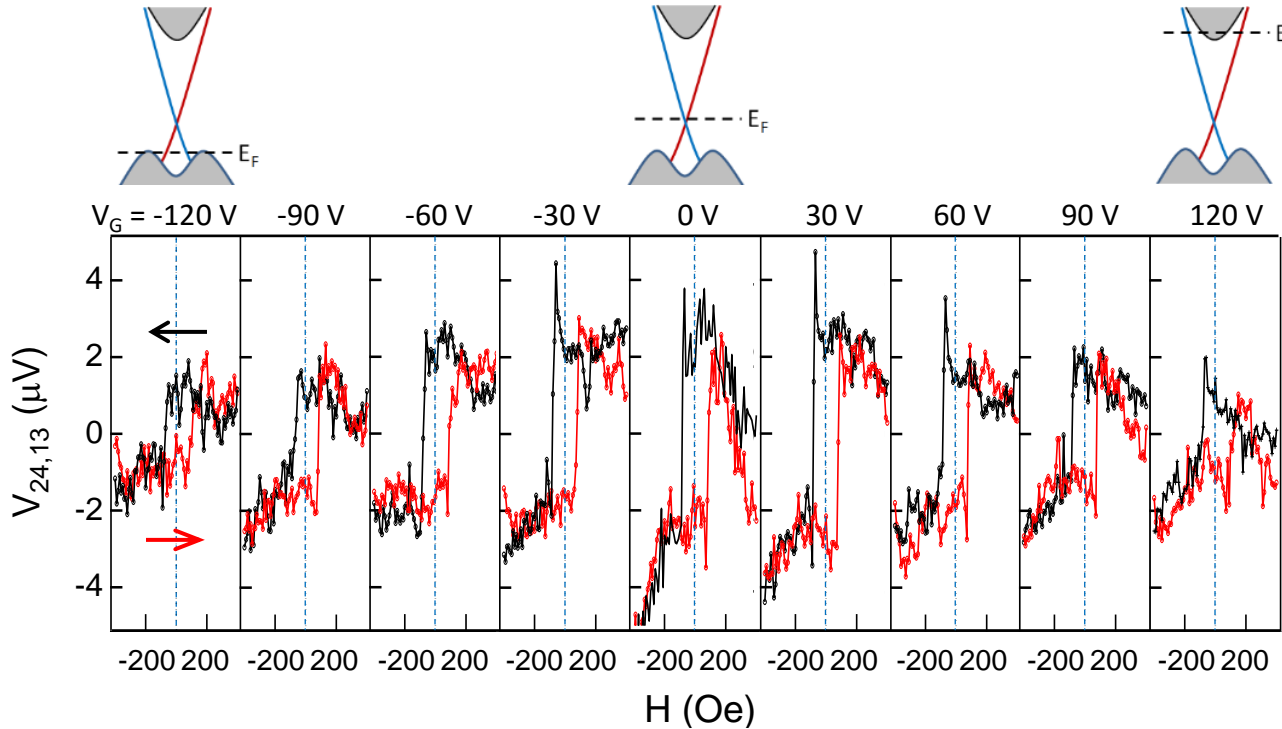
- Larger spin signals with the chemical potential near DP.
→ Minimal bulk effects.
- No change in the direction of spin polarization across DP.

Lee *et al.*, *Phys. Rev. B* **92**, 155312 (2015)



Detection of spin polarization with back gating

Lee *et al.*, *Phys. Rev. B*
92, 155312 (2015)



Spin polarization of TI

$$\Delta V = IR_B P_{TI} P_{FM}$$

- ΔV : detected spin signal
- I : bias current
- P_{FM} : spin polarization of Py
- R_B : ballistic resistance

For μ near Dirac point,

$$P_{TI} = \mathbf{0.42 \pm 0.15} \text{ (Device A), } \mathbf{0.78 \pm 0.26} \text{ (Device B)}$$

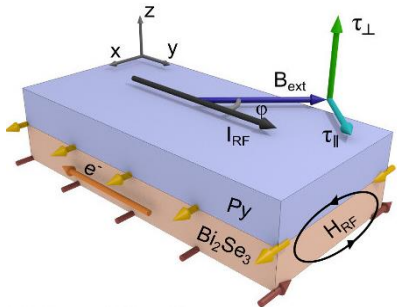
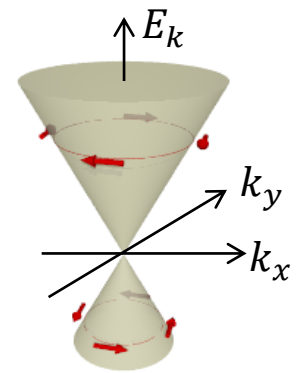
→ Comparable with theoretical calculations
 $\mathbf{0.5 \sim 2/\pi}$.

Hong *et al.*, *Phys. Rev. B*. **86**, 085131 (2012)

Yazyev *et al.*, *Phys. Rev. Lett.* **105**, 266806 (2010)

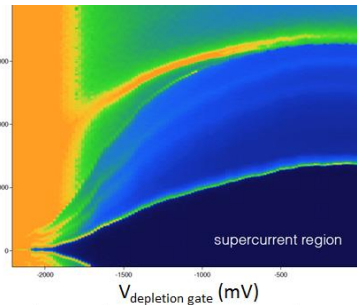
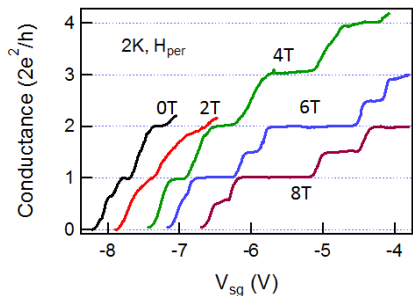
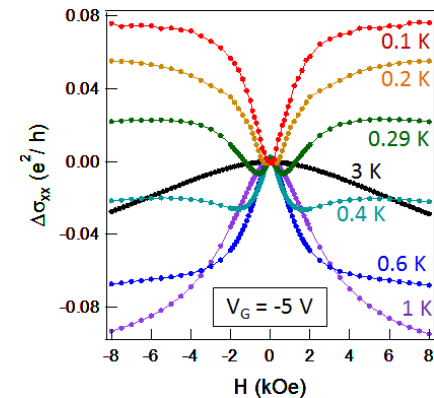
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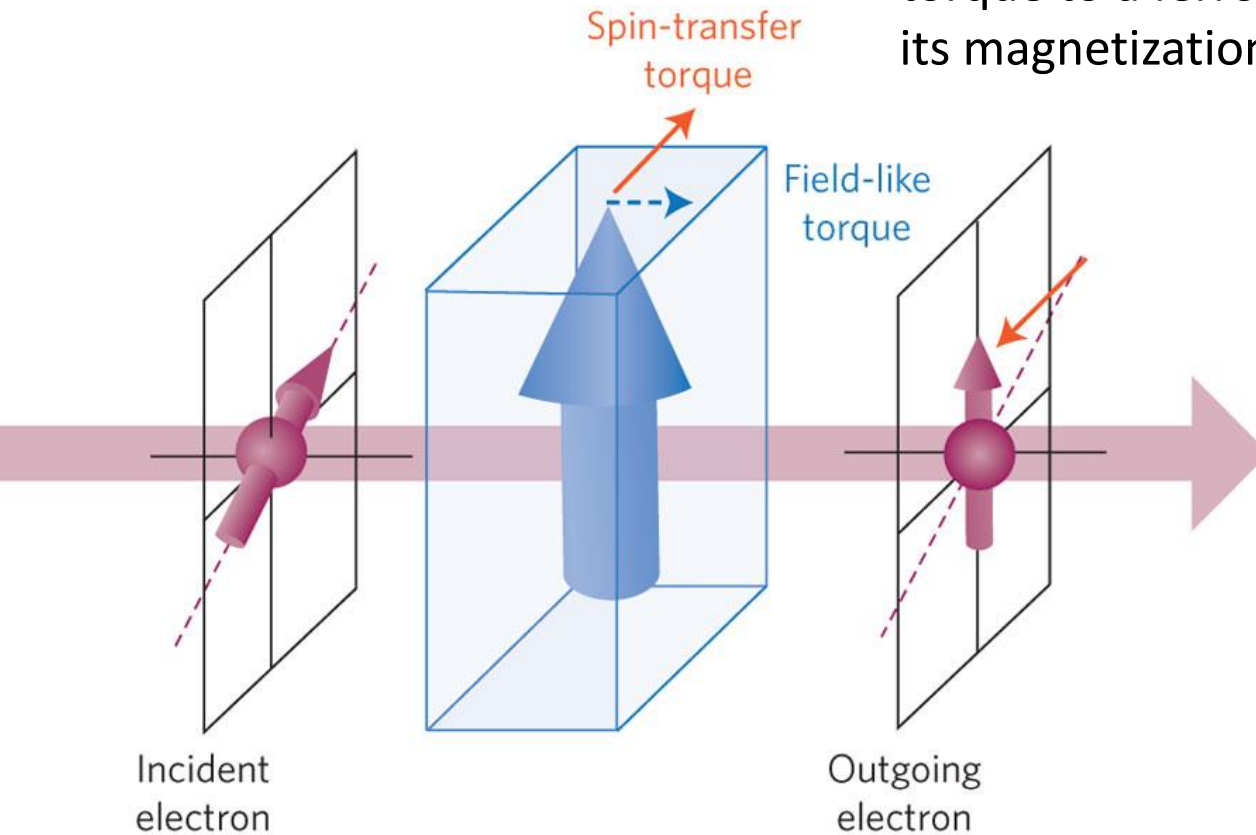
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Spin-transfer torque

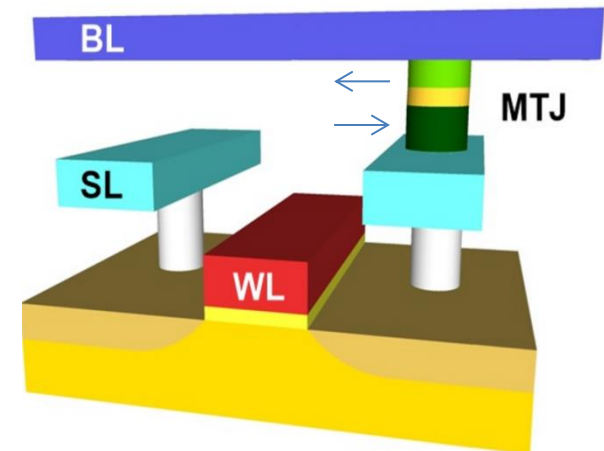
Incident spin-polarized electron exerts a torque to a ferromagnet to precess/reverse its magnetization orientation.



Brataas, *et al.*, *Nat. Mater.* **11**, 372 (2012)

Applications:

- Magnetic random access memory (MRAM)
- Written by spin-transfer torque, read by tunnel magnetoresistance.

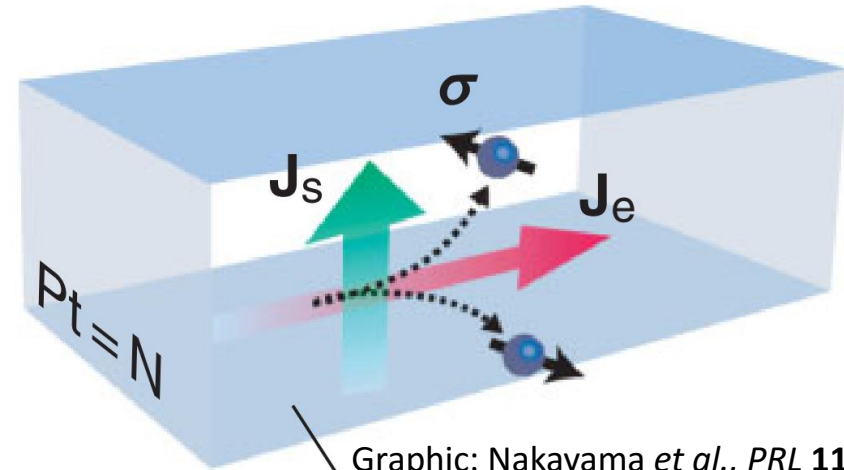


cspin.umn.edu/newsletter

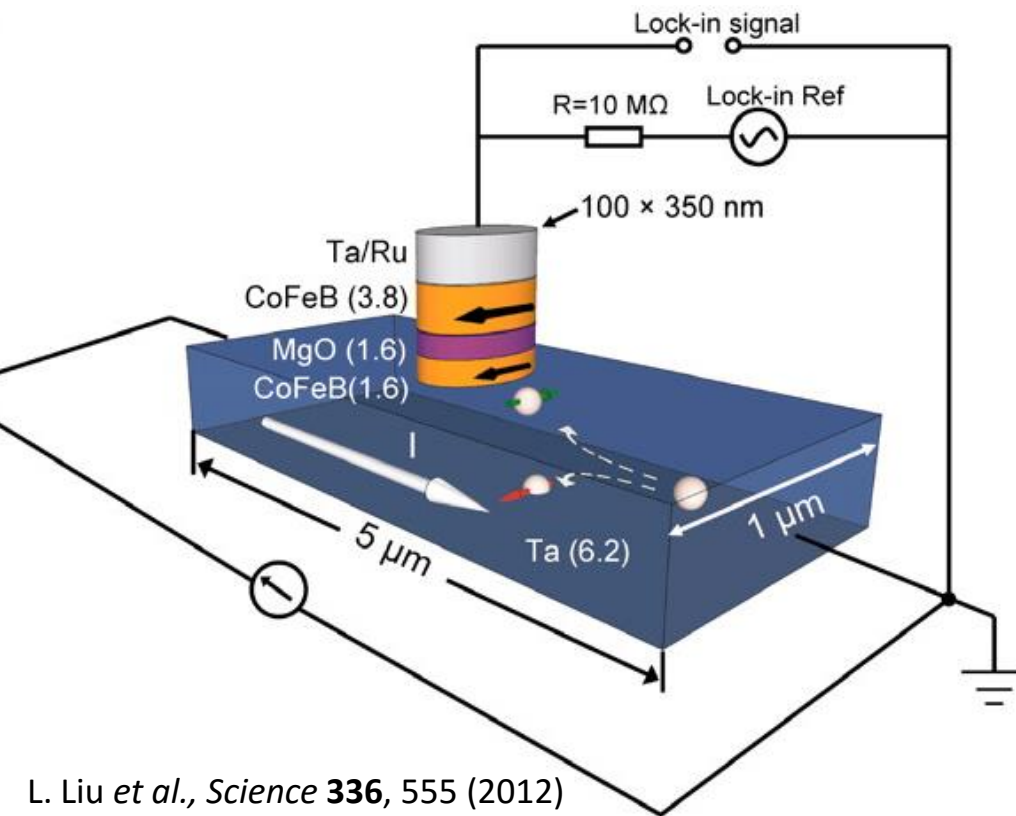
Source of spin current for spin-transfer torque

Spin Hall effect:

Strong spin-orbit interaction induces spin polarization and spin accumulation in semiconductors and heavy metals.

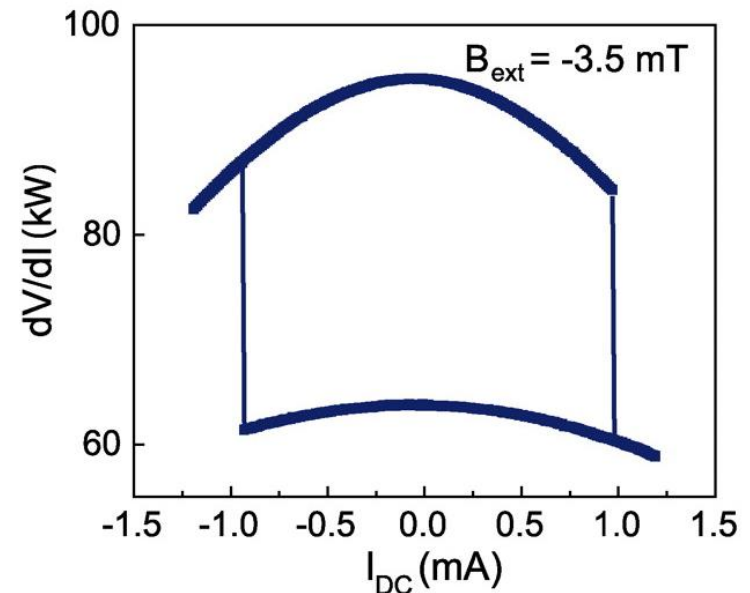


Graphic: Nakayama *et al.*, *PRL* **110**, 206601 (2013)



L. Liu *et al.*, *Science* **336**, 555 (2012)

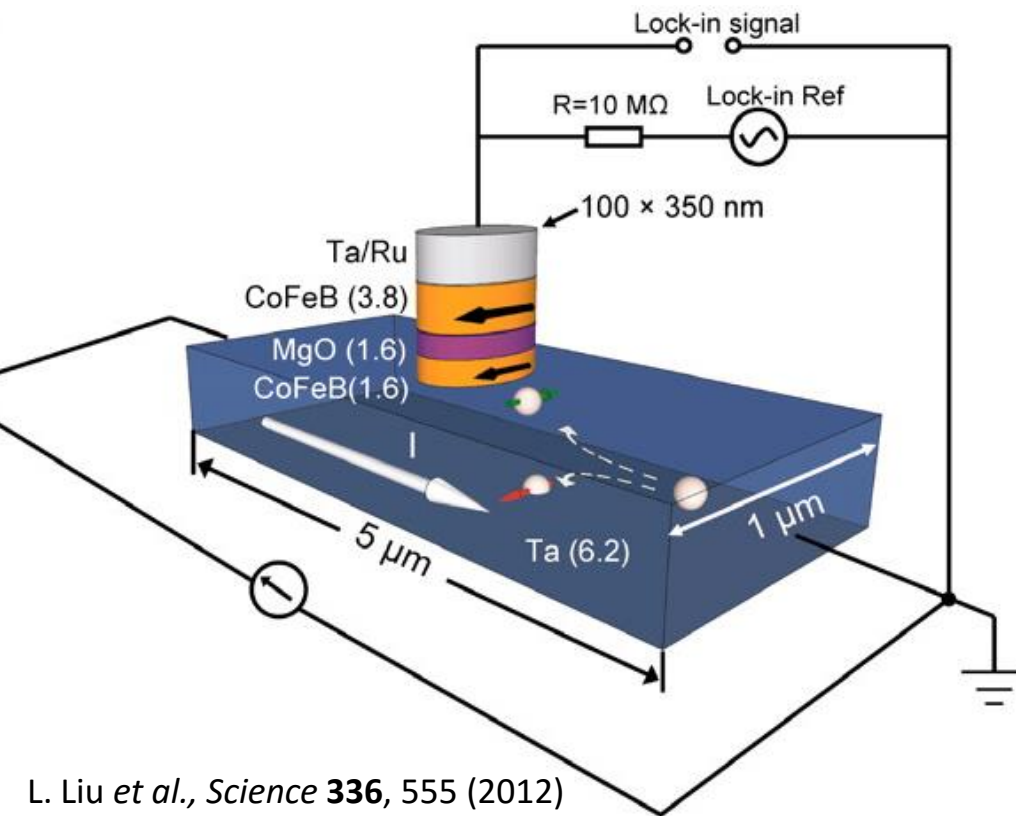
- Spin-transfer torque by spin Hall effect
- Measured by tunnel magnetoresistance



Source of spin current for spin-transfer torque

Spin Hall effect:

Strong spin-orbit interaction induces spin polarization and spin accumulation in semiconductors and heavy metals.



L. Liu *et al.*, *Science* **336**, 555 (2012)

- Spin-transfer torque by spin Hall effect
- Measured by tunnel magnetoresistance

Charge-to-spin conversion efficiency:

- Spin Hall angle
- Spin torque ratio

$$\theta_{||} \equiv \frac{2e}{\hbar} \frac{J_{s,||}}{J} = \frac{2e}{\hbar} \frac{\sigma_{s,||}}{\sigma}$$

Currently studied materials

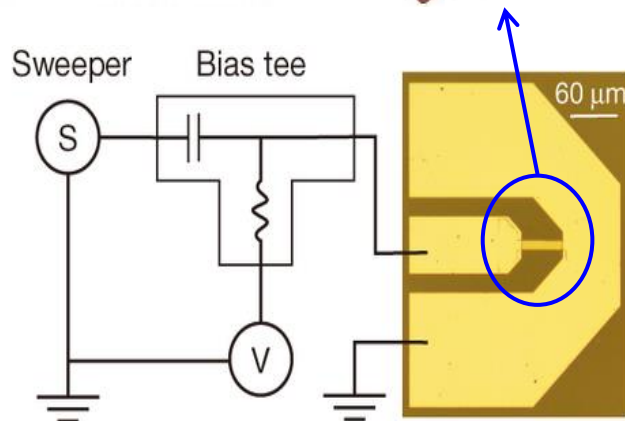
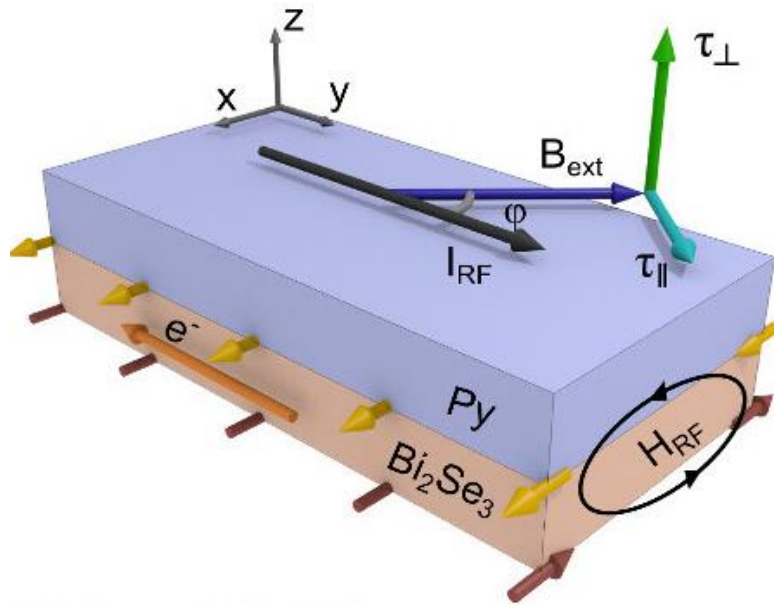
Pt	β -Ta	Cu(Bi)	β -W
0.08	0.15	0.24	0.3

Liu *et al.*, *PRL* **106**, 036601 (2011)
 Miron *et al.*, *Nature* **476**, 189 (2011)
 Niimi *et al.*, *PRL* **109**, 156602 (2012)
 Pai *et al.*, *APL* **122**, 101404 (2012)

More efficient charge-to-spin conversion by topological insulators?

Spin-transfer torque by a topological insulator

Current-induced spin current in Bi_2Se_3 exerts strong spin-transfer torque to Py layer.



Spin-torque measured by ferromagnetic resonance (FMR)

rf-current through $\text{Bi}_2\text{Se}_3/\text{Py}$ with B_{ext}



Precessing Py magnetization & oscillating TI spin current

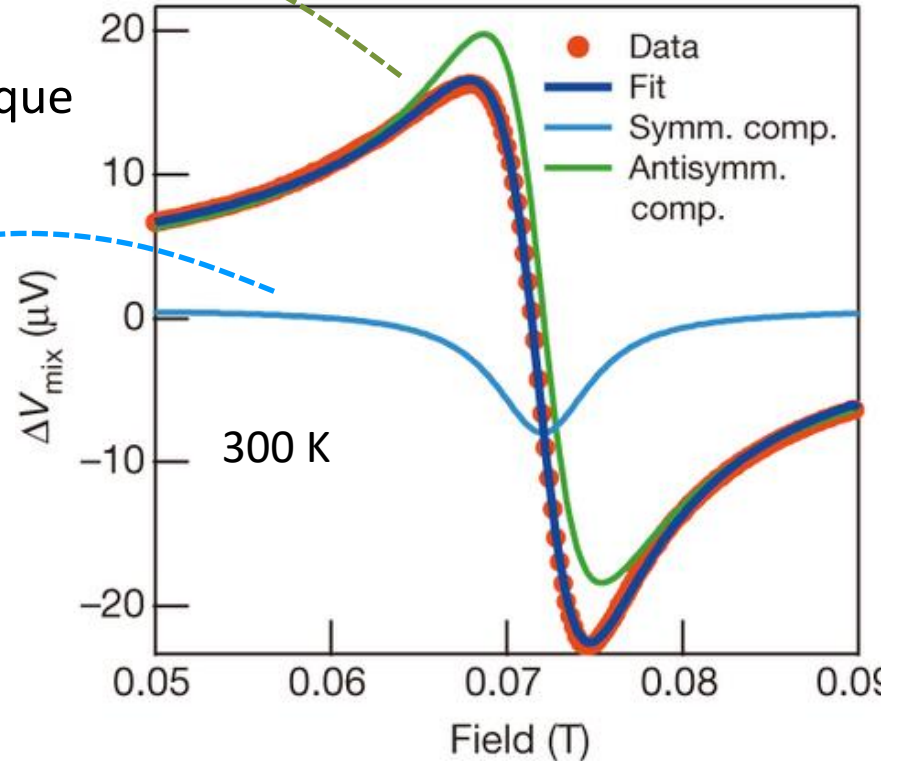
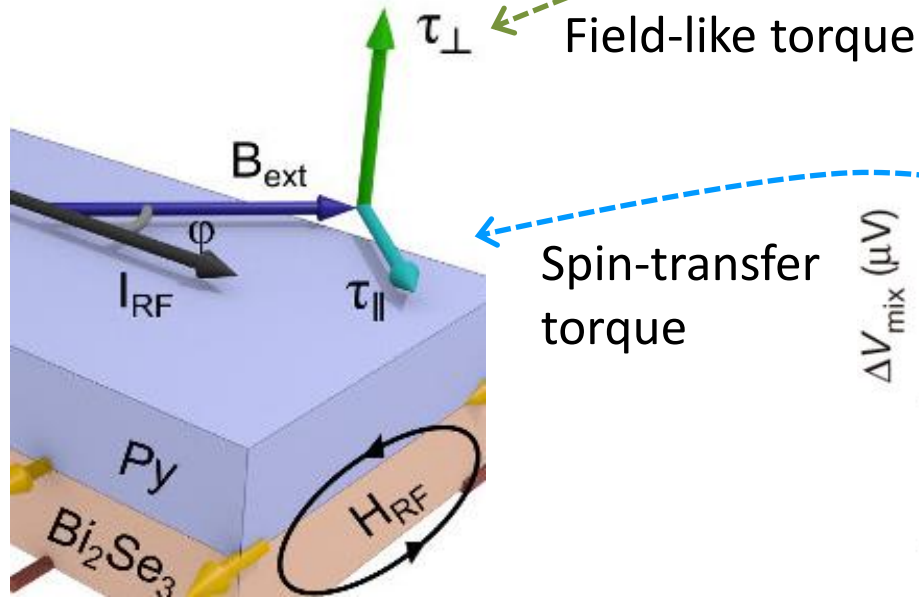


Precession perturbed by torques from the spin current



Measure DC mixing voltage

Spin-transfer torque by a topological insulator



Spin-transfer torque component could be induced by TI surface states or Rashba states.

→ sign consistent with topological surface states

$$\theta_{\parallel} \equiv \frac{2e}{\hbar} \frac{J_{s,\parallel}}{J}$$

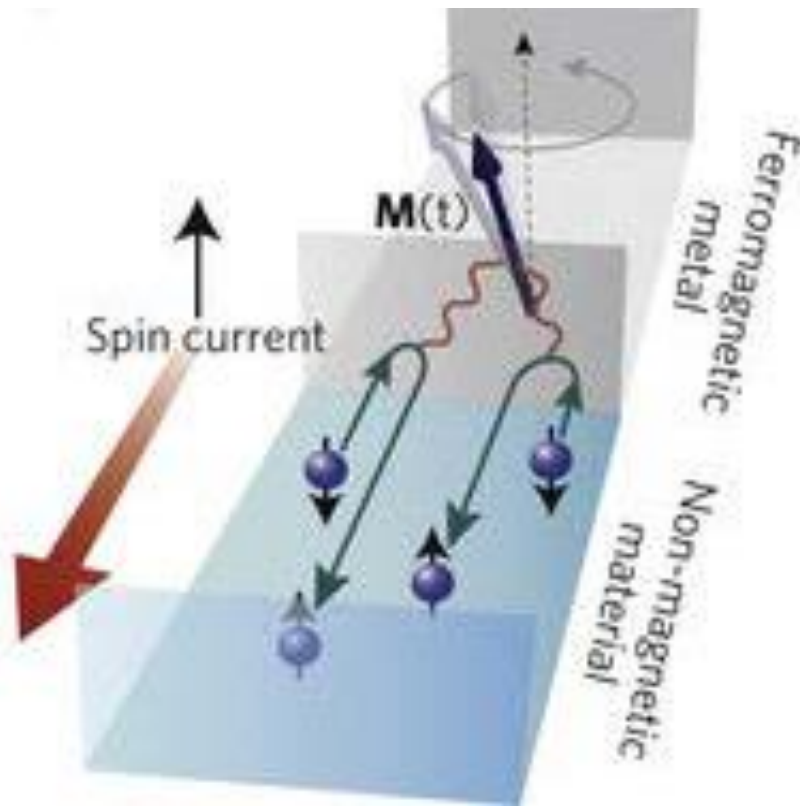
Bi ₂ Se ₃	Pt	β -Ta	Cu(Bi)	β -W
1 ± 0.4	0.08	0.15	0.24	0.3

Efficient charge-to-spin conversion by topological surface states.

Spin-to-charge conversion by spin pumping

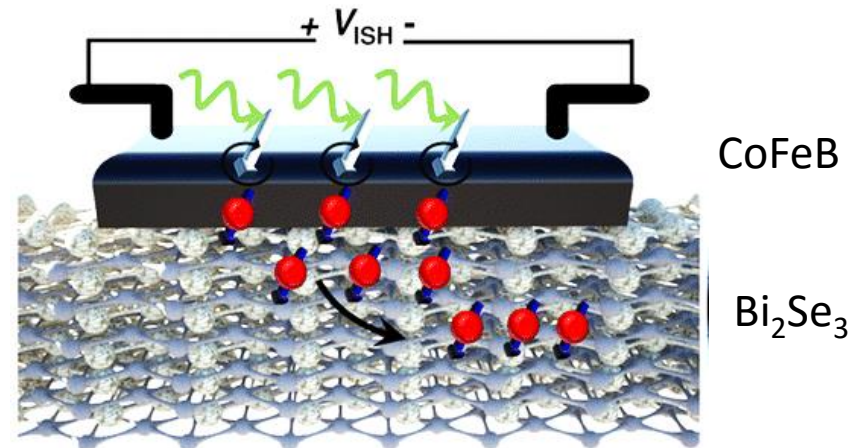
Spin pumping:

Precessional magnetization dynamics of a ferromagnet injects a spin current into an adjacent layer.



Ando *et al.*, *Nature Mater.* **10**, 655 (2011)

Spin pumping into Bi_2Se_3 from CoFeB



FMR-driven spin pumping

Precessing FM magnetization by external rf-magnetic field



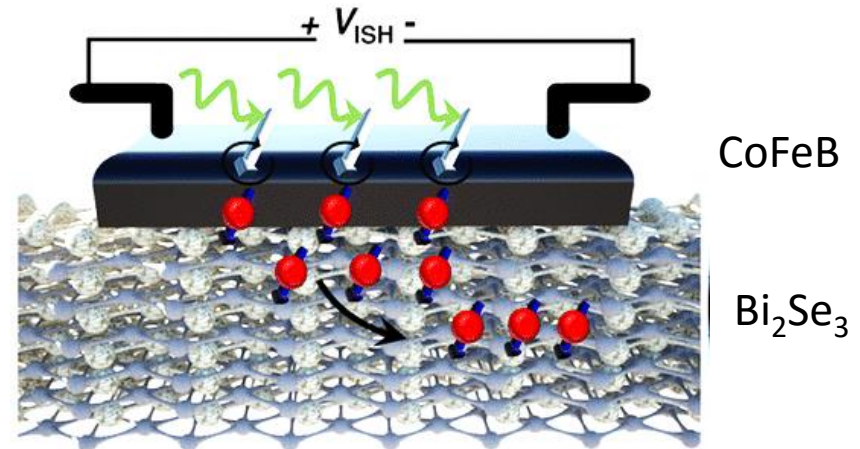
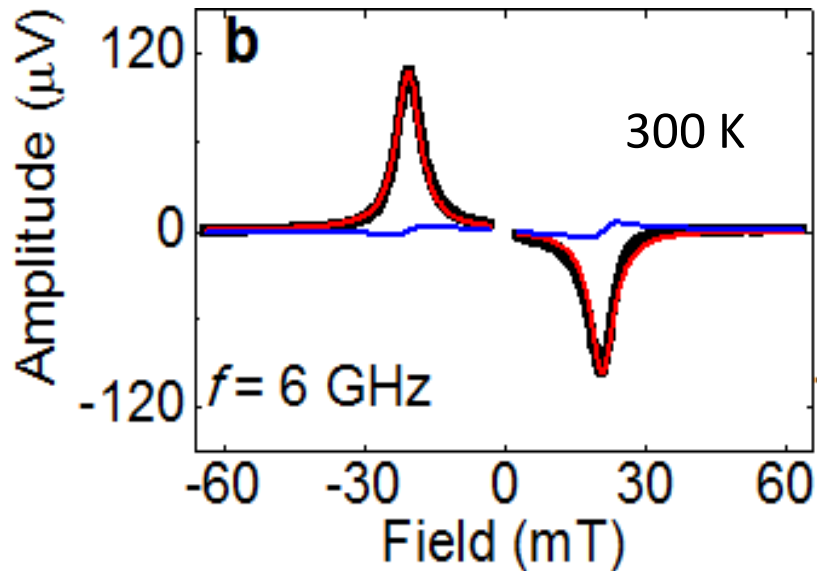
Injecting a spin current into Bi_2Se_3



Spin current into charge current by inverse spin Hall effect

Jamali, Lee *et al.*, *Nano Lett.* **15**, 7126 (2015)

Spin pumping into Bi_2Se_3 from a metallic FM



Jamali, Lee *et al.*, *Nano Lett.* **15**, 7126 (2015)

$$\theta_{\parallel} \equiv \frac{2e}{\hbar} \frac{J_{s,\parallel}}{J} = \frac{2e}{\hbar} \frac{\sigma_{s,\parallel}}{\sigma}$$

Resulting spin Hall angle:
0.43 at room temperature

Spin-to-charge conversion mechanisms

- 1) Inverse Edelstein effect (IEE) by surface states
- 2) Inverse spin Hall effect (ISHE) by bulk

→ dominant!

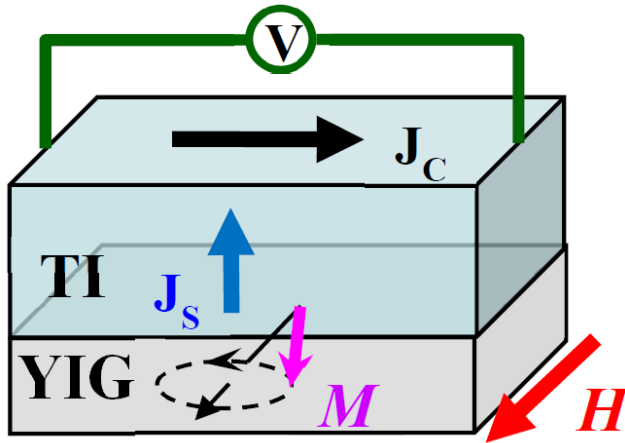
Issues in spin-charge conversion experiments

- “Surface effect” vs “Bulk effect”
- Shunting of charge current by metallic FM

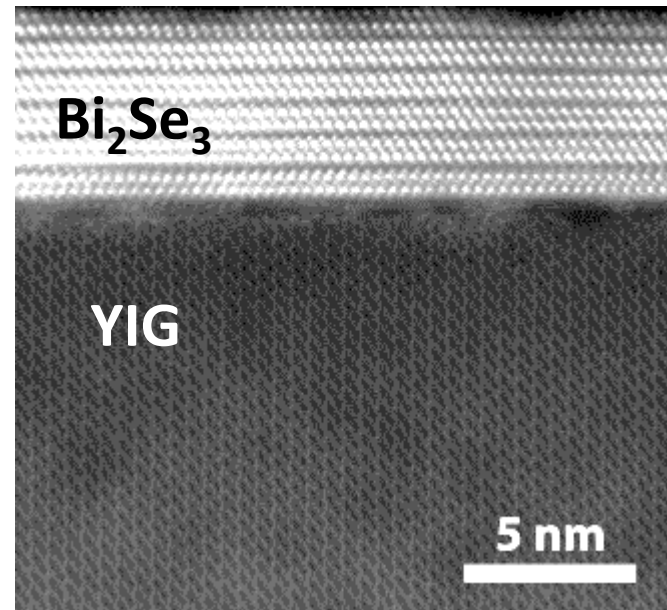
Spin pumping into Bi_2Se_3 from an insulating FM

Avoiding shunting of current using an Insulating ferromagnet.

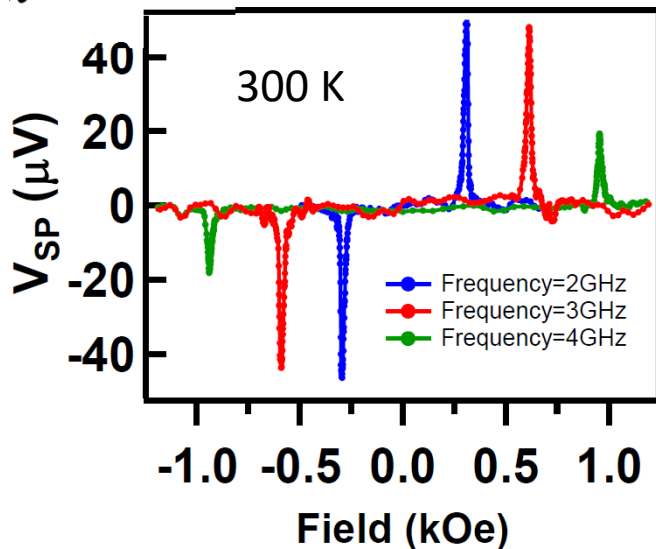
$\text{Y}_3\text{Fe}_5\text{O}_{12}$ (YIG) : insulating ferromagnet with a low damping constant



MBE-grown Bi_2Se_3 on YIG



FMR-driven spin pumping

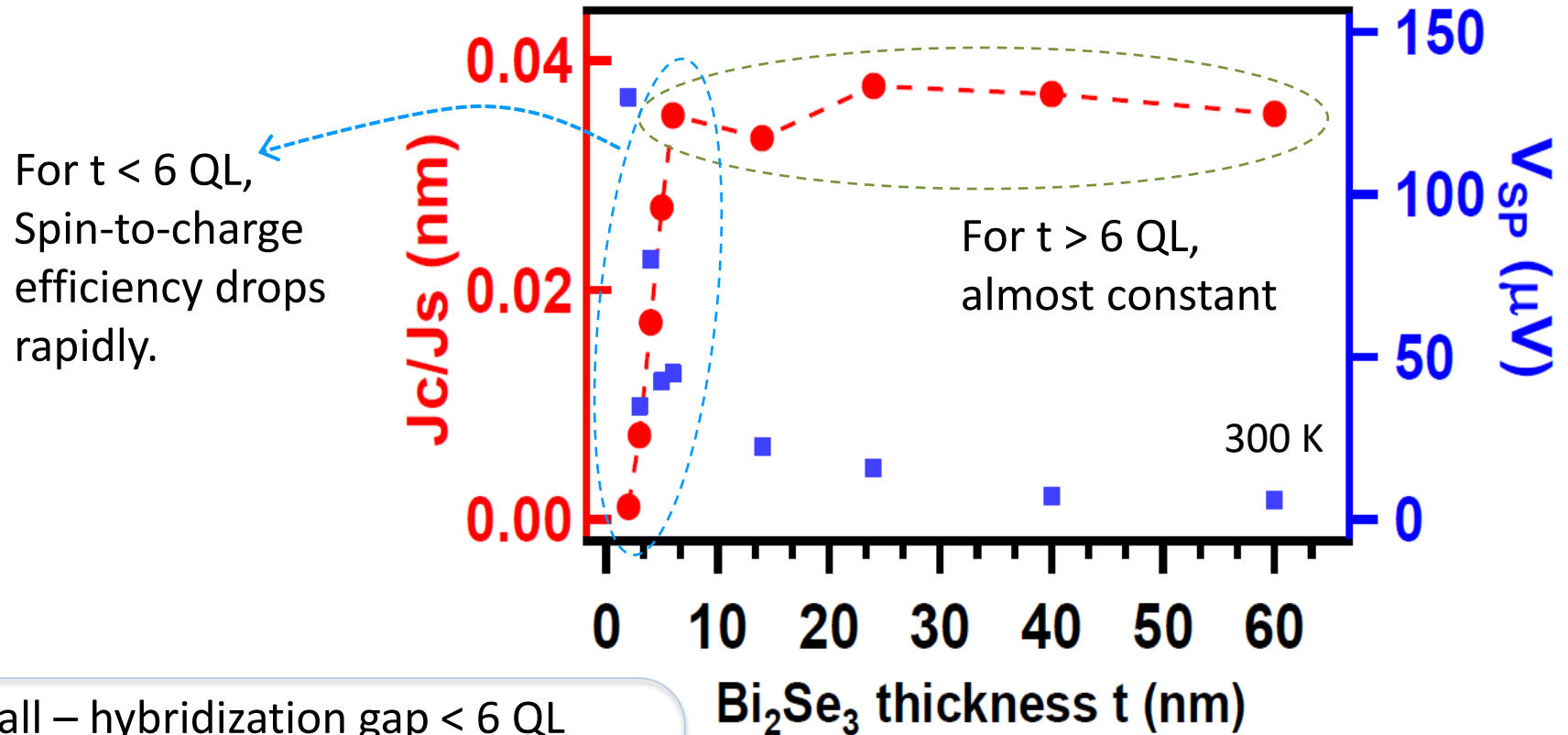


- Two-step growth of Bi_2Se_3 film
- Disorders at the interface (~ 1 nm)

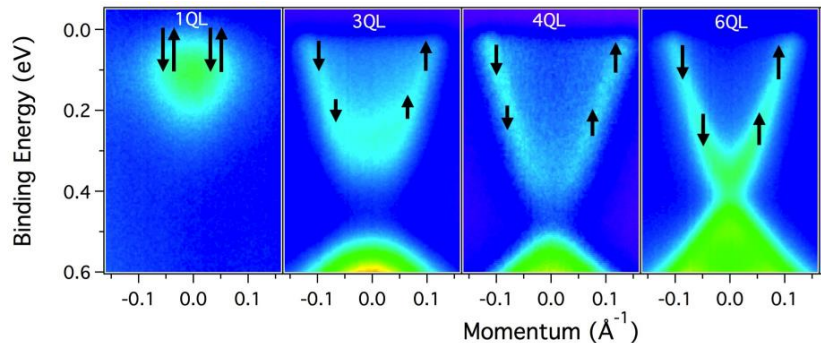
Wang, Kally, Lee *et al.*, Phys. Rev. Lett. **117**, 076601 (2016)

Thickness dependence of spin pumping using YIG

J_c/J_s = Spin-to-charge conversion efficiency. (J_c , J_s : spin, charge current density)



Recall – hybridization gap < 6 QL



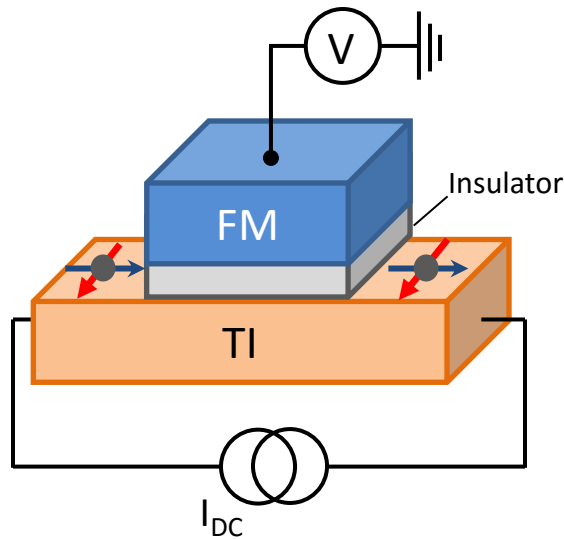
M. Neupane *et al.*, *Nat. Comm.* **5**, 3841 (2014)

Surface states are dominant in spin-to-charge conversion in Bi_2Se_3 .

Wang, Kally, Lee *et al.*, *Phys. Rev. Lett.* **117**, 076601 (2016)

Summary of Part 1 and 2

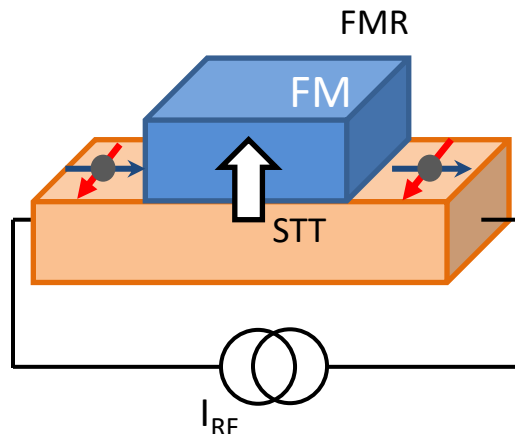
Detecting Spin polarization in topological surface states



- DC current through TI generates spin polarization.
- Use FM as a spin detector.

Lee *et al.*, *PRB* (2015)

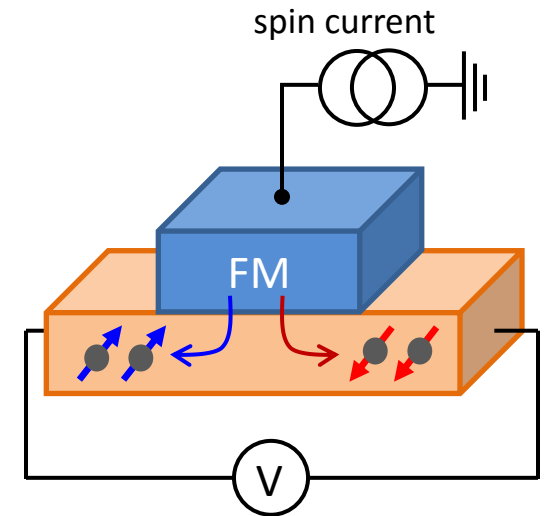
Charge-to-spin conversion Spin-transfer torque by TI



- rf current through TI exerts spin-transfer torque onto FM.
- Detected by FMR.

Mellnik, Lee *et al.*, *Nature* (2014)

Spin-to-charge conversion Spin pumping into TI



- Spin current from FM into TI.
- Detection of electromotive force by spin-orbit interaction.

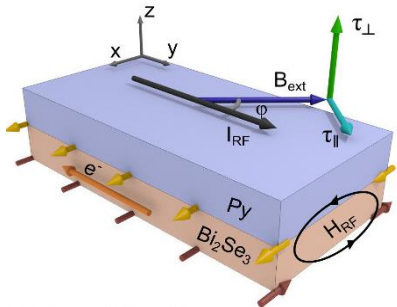
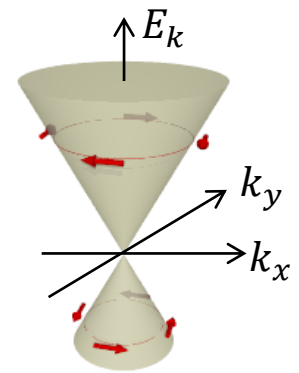
Jamali, Lee *et al.*, *Nano Lett.* (2015)

Wang, Kally, Lee *et al.*, *PRL* (2016)

Anything more to consider near Dirac point?

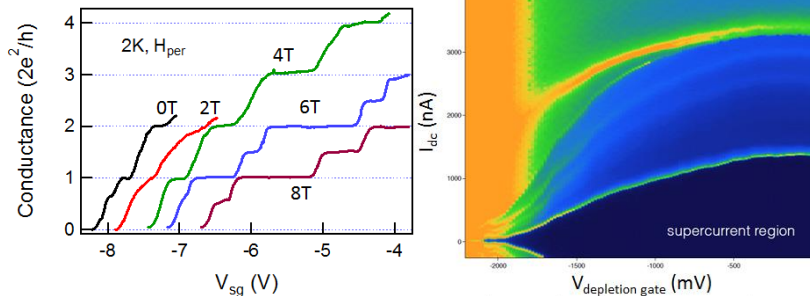
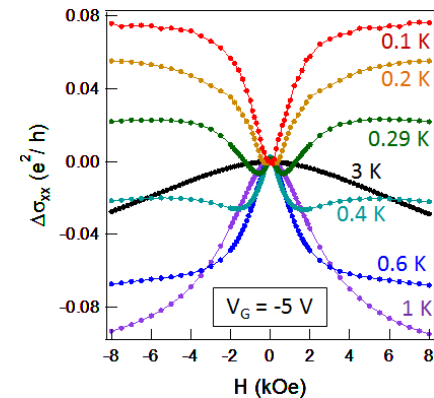
Outline

- Spin-polarized surface states in 3D topological insulators
 - Electrical detection of spin polarization



- Spin-charge conversion in topological insulators
 - Charge-to-spin conversion: spin transfer torque
 - Spin-to-charge conversion: spin pumping

- Modified surface states by breaking time-reversal symmetry
 - Interfacing a ferromagnetic insulator to a TI



- Topological superconductivity towards quantum computation
 - 1D transport and superconducting proximity effect in epi-Al/2DEG

Magnetic surface phases modified by broken TR symmetry

TI surface states can be modified by introducing magnetism.

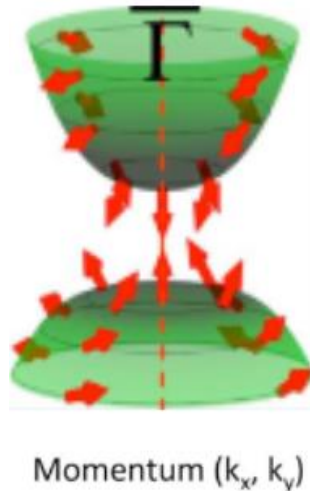
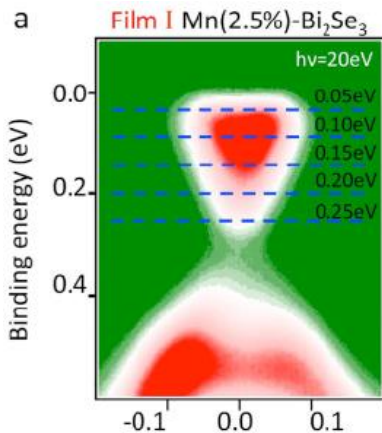
$$H_{surf}(k_x, k_y) = \hbar v_F(\sigma^x k_y - \sigma^y k_x) + \sum_{a=x,y,z} m_a \sigma^a$$

$$E_k = \pm \sqrt{(\hbar v_F k_y + m_x)^2 + (\hbar v_F k_x - m_y)^2 + m_z^2}$$

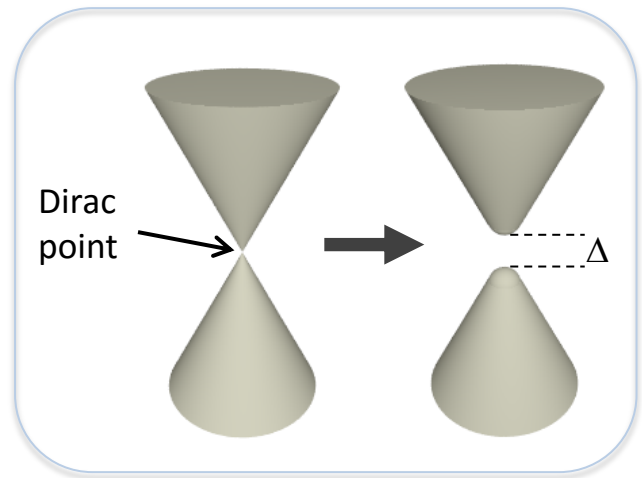
m_z induces an energy gap.

Hedgehog-like spin texture

Mn-Bi₂Se₃



Xu *et al.*, *Nat. Phys.* **8**, 616 (2012)



- m_z can be introduced by
 - magnetic doping
 - interfacing a magnetic layer

Interfacing insulating FM to a topological insulator

Interfacing a magnetic layer breaks time-reversal symmetry and opens a gap at the Dirac point.

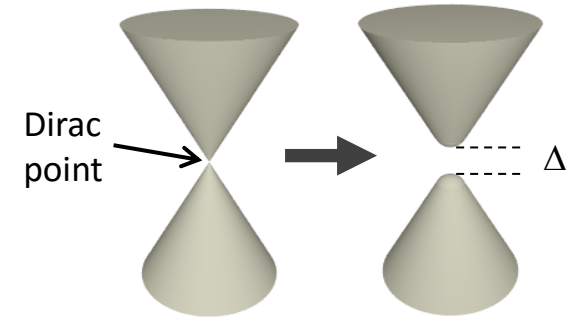
Experimental challenge:

How to measure the modified surface state from the buried interface?

- No ARPES and no STM. Probably electrical transport?

Need to develop...

- (1) Topological insulator with E_F near the Dirac point.
- (2) Insulating FM with perpendicular magnetization.
- (3) Clean, well-defined TI/FM interface.

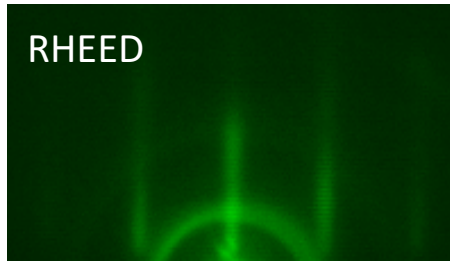


Topological insulator/GaMnAs heterostructure by MBE

- Topological insulator with E_F near the Dirac point.
- Insulating FM with perpendicular magnetization.
- Clean, well-defined TI/FM interface.

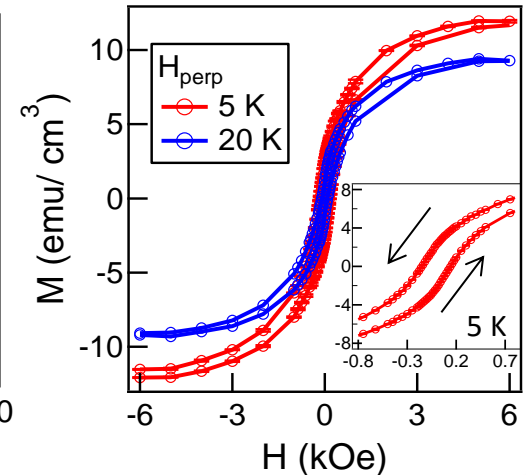
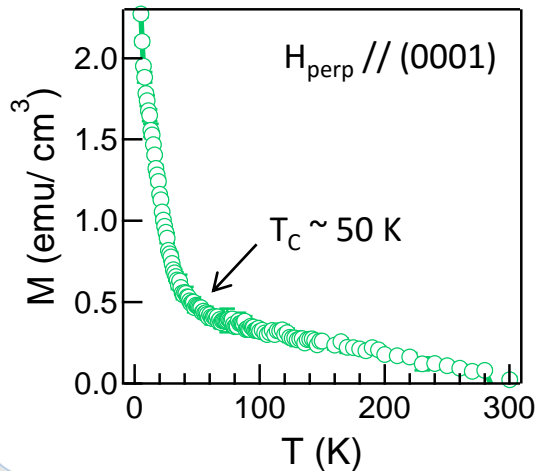
[TI] $(\text{Bi,Sb})_2(\text{Te,Se})_3$

- ~8 QL TI thin film grown by MBE.
- E_F in the surface state.
(Bi:Sb \approx 1:1, Te:Se \approx 2:1)
- Further E_F tuning by top-gating.



[FM] $(\text{Ga}_{1-x}\text{Mn}_x)\text{As}$ on InP (111)A

- Perpendicular component of magnetization.
- Low Mn doping ($\text{Mn}\approx 0.05$) \rightarrow highly resistive.

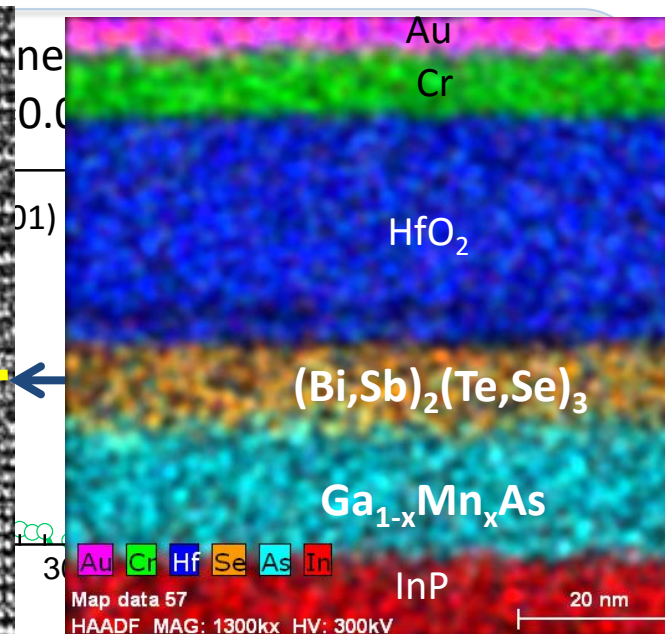
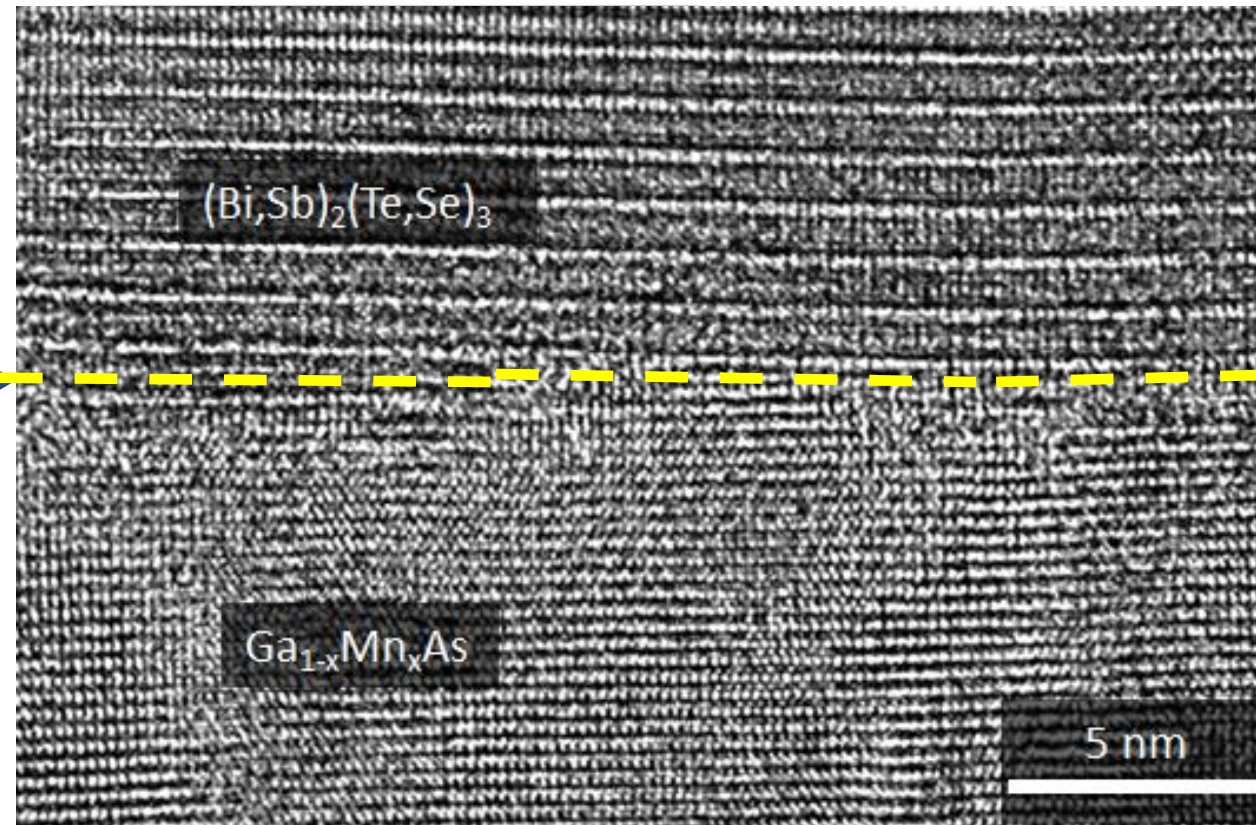


Topological insulator/GaMnAs heterostructure by MBE

- ✓ Topological insulator with E_F near the Dirac point.
- ✓ Insulating FM with perpendicular magnetization.
- ✓ Clean, well-defined TI/FM interface.

[TI] $(\text{Bi,Sb})_2(\text{Te,Se})_3$

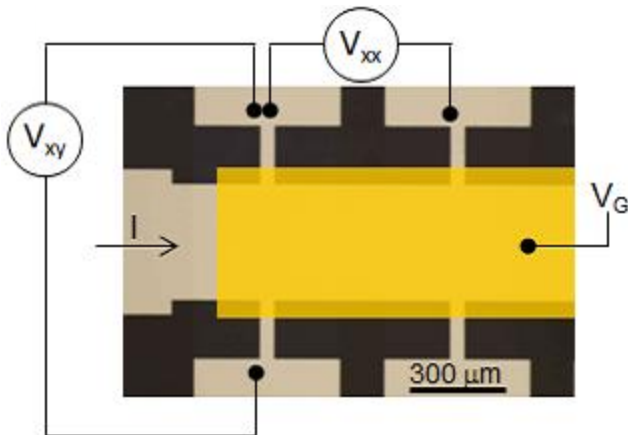
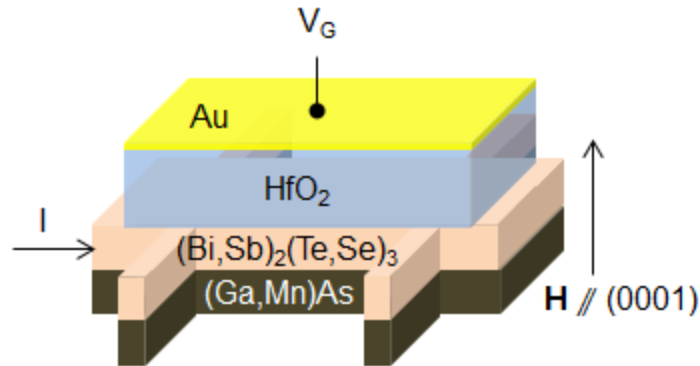
[FM] $(\text{Ga}_{1-x}\text{Mn}_x)\text{As}$ on InP (111)A



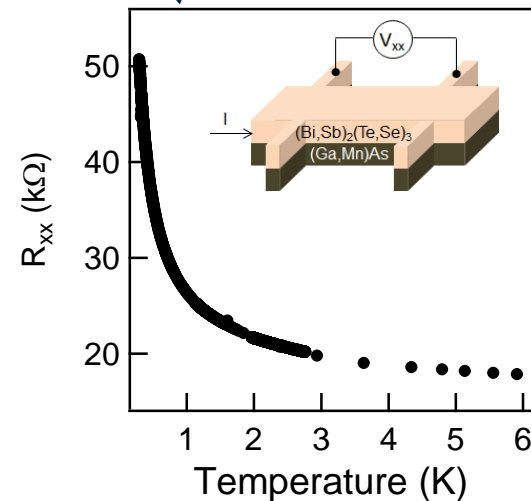
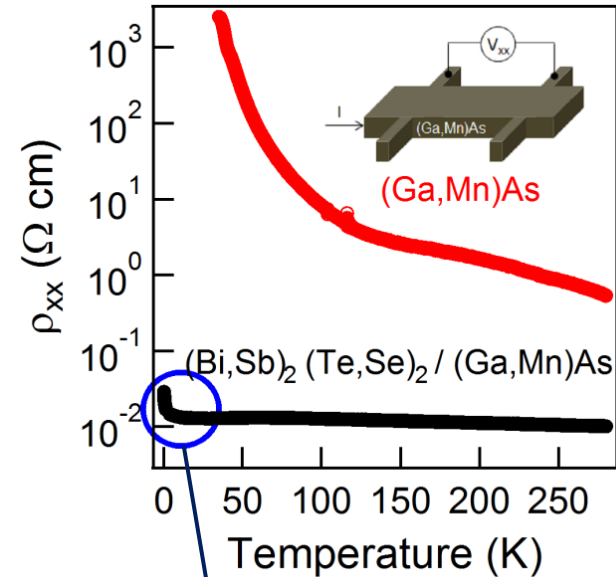
Energy dispersive x-ray spectroscopy

Electrical transport using a top-gated Hall bar device

Our approach: Transport measurements for evidencing the modified surface states



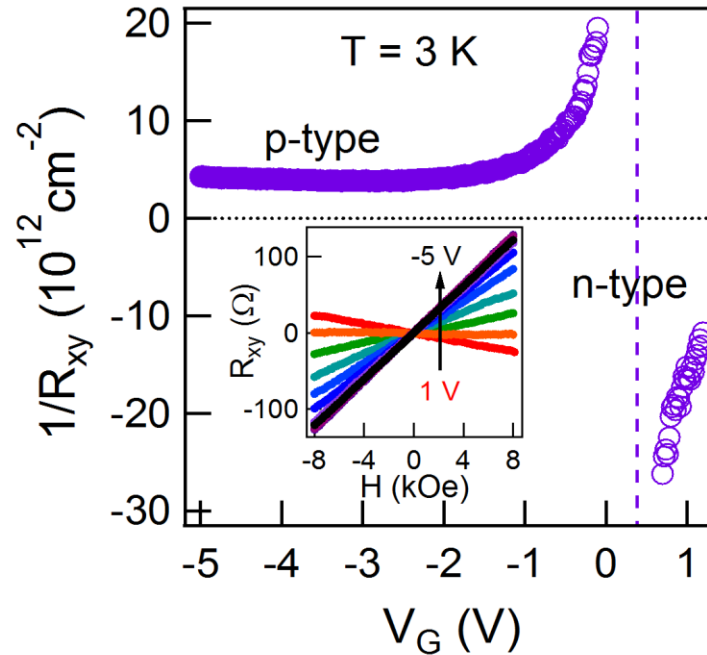
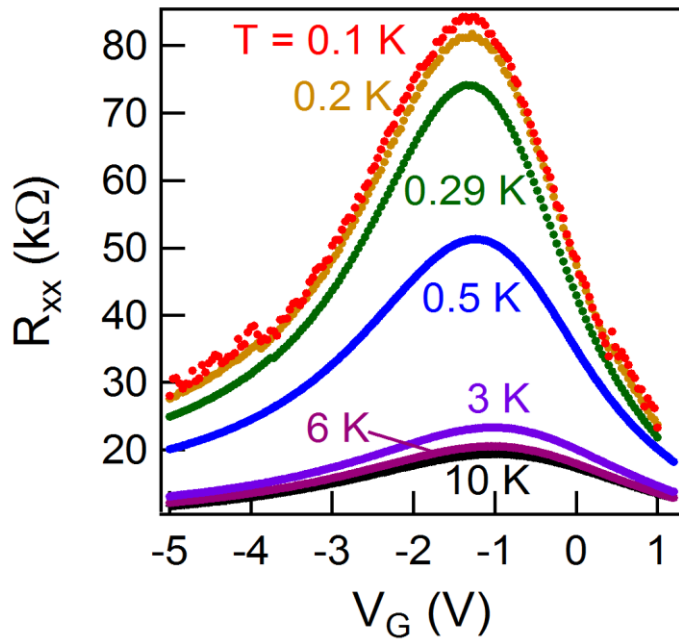
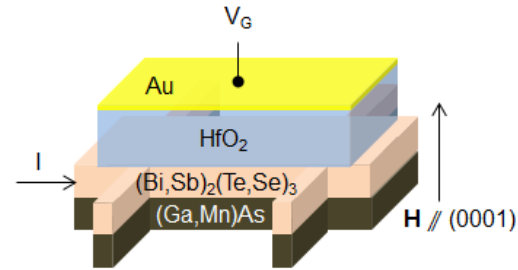
- HfO₂ (25 ~ 30 nm) by evaporation and ALD.
- Hall bar: 650 μm x 400 μm
- Patterned by standard photo-lithography.



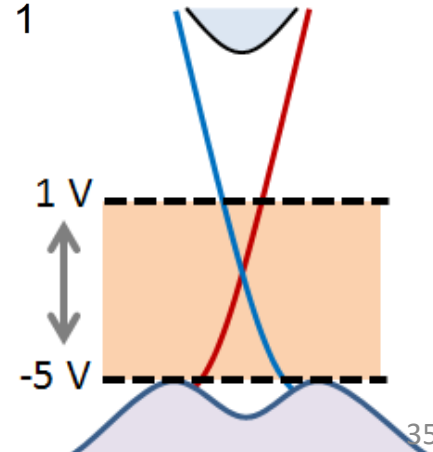
→ Current flows through TI layer.

Ambipolar transport in $(\text{Bi,Se})_2(\text{Se,Te})_3/\text{GaMnAs}$

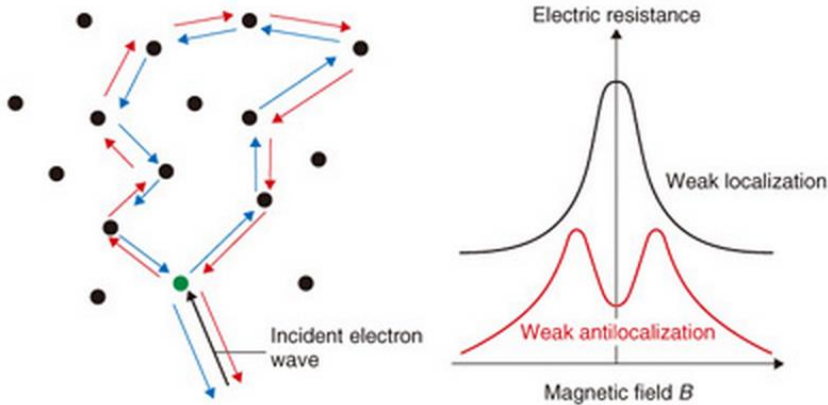
- Gate-voltage dependence reveals ambipolar transport.



- n-type to p-type change.
- In the V_G range, chemical potential is in a region where bulk is depleted.



Quantum corrections of weak (anti)localization



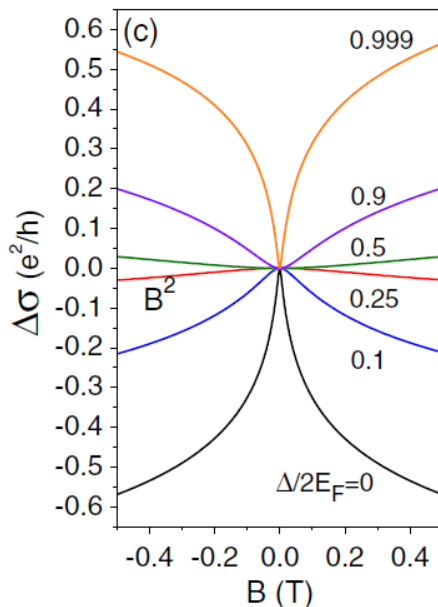
<https://www.nttreview.jp/anqtest/archive/ntttechnical.php?contents=ntr201209fa4.html>

Weak (anti-)localization

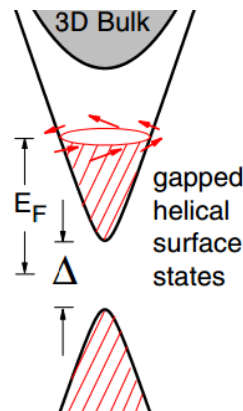
Constructive (destructive) interference between time-reversal paths around self-intersecting loops when phase coherence is maintained leads to WL (WAL).

- [3D topological insulators: WAL](#)
- from the π Berry phase of electrons going around the Fermi circle of the Dirac surface state.

Crossover between WAL and WL in TI



WL
↑
↓
WAL



Modified Berry phase

$$\psi_{\mathbf{k}}(\mathbf{r}) = \pi \left(1 - \frac{\Delta}{2E_F} \right)$$

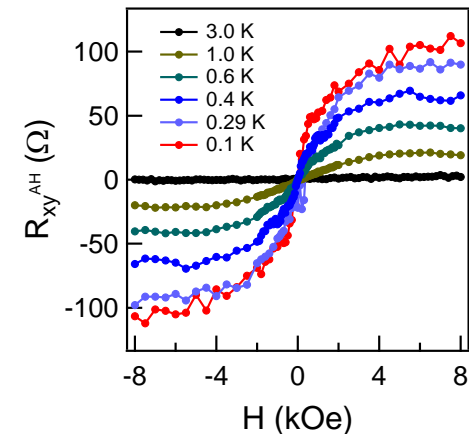
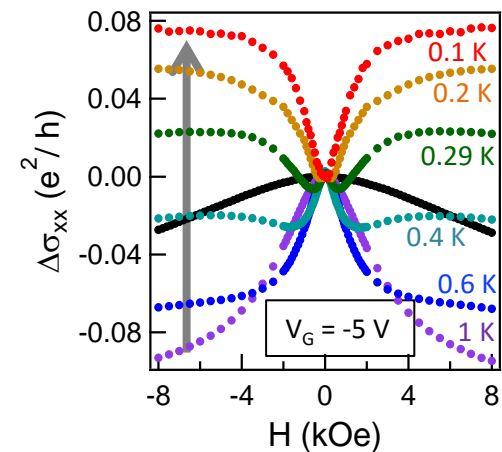
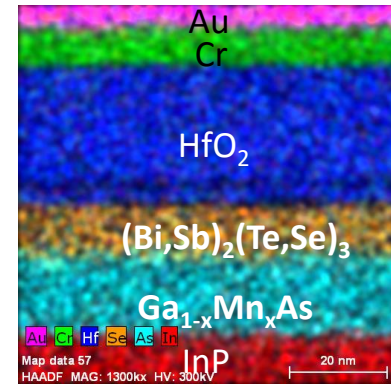
(Δ : gap, E_F : Fermi energy)

$$\left\{ \begin{array}{l} \text{For } \Delta = 0 \text{ or } E_F \gg \Delta : \varphi = \pi \\ \text{For } \Delta = 2E_F : \varphi = 0 \end{array} \right.$$

Lu *et al.*, Phys. Rev. Lett. **107**, 076801 (2011).

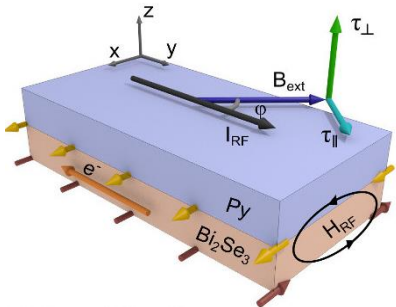
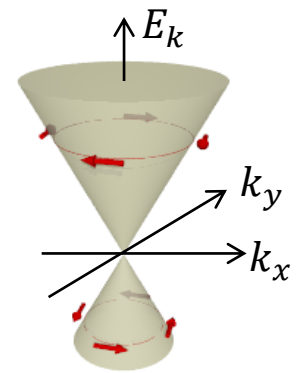
Summary of Part 3

- To evidence the modified surface states by broken TR symmetry, we developed $(\text{Bi,Sb})_2(\text{Te,Se})_3/(\text{Ga,Mn})\text{As}$ heterostructure.
- Crossover between WAL and WL as well as systematic emergence of AHE with varying T and V_G .
→ interpreted as a result of the Berry phase modification by a gap opening and chemical potential tuning.
- The results suggest that systematic changes in WAL and AHE can be used as probes of modified surface states.



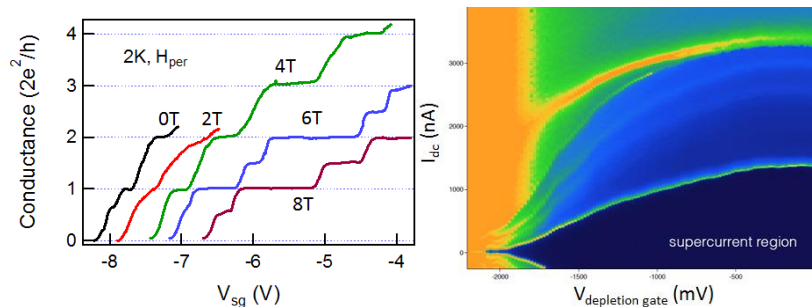
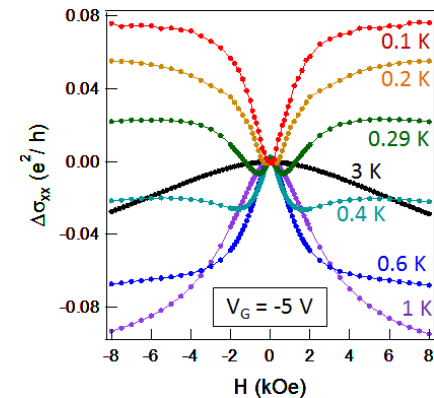
Outline

1. Spin-polarized surface states in 3D topological insulators
 - Electrical detection of spin polarization



2. Spin-charge conversion in topological insulators
 - Charge-to-spin conversion: spin transfer torque
 - Spin-to-charge conversion: spin pumping

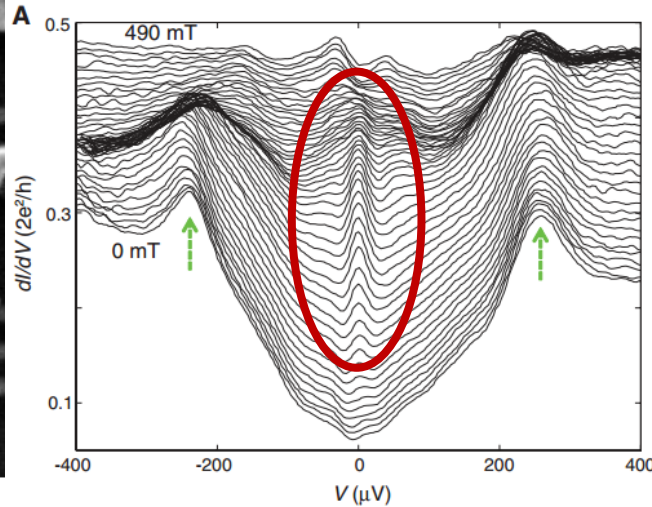
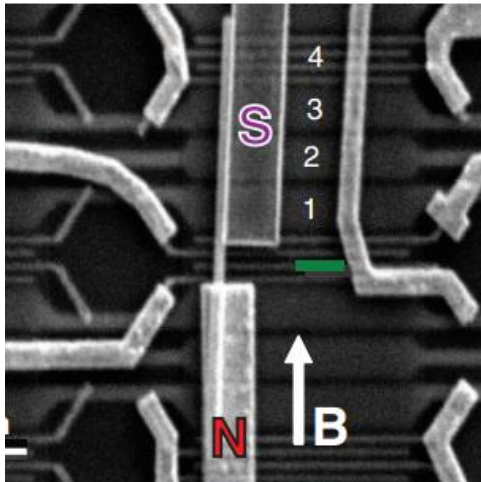
3. Modified surface states by breaking time-reversal symmetry
 - Interfacing a ferromagnetic insulator to a TI



4. Topological superconductivity towards quantum computation
 - 1D transport and superconducting proximity effect in epi-Al/2DEG

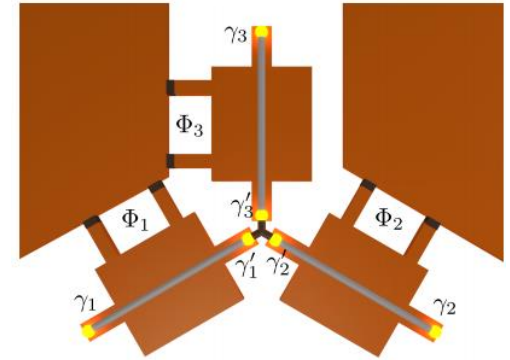
Motivation for epi-Al/InAs QW

Majorana fermions – zero energy mode



Mourik *et al.*, *Science* **336**, 1003 (2012)

Braiding Majorana fermions for quantum computation



Beenakker, *Annu. Rev. Condens. Matter* **4**, 113 (2013)

Current issues: 1) soft gap of superconducting proximity, 2) difficulties in large scale fabrication using nanowires.

Current approach

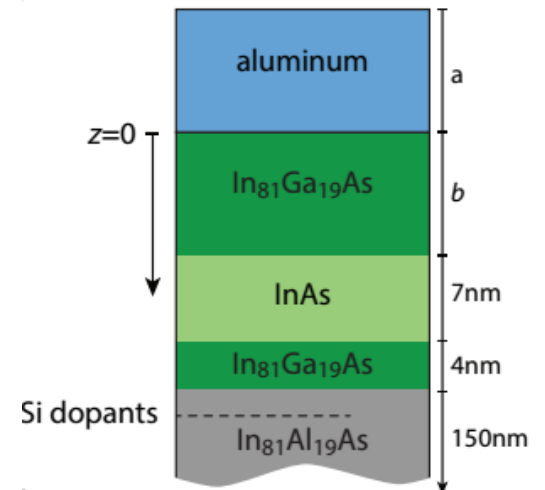
- 1D nanowire with strong spin-orbit coupling
- Superconducting proximity <sputtered vs epitaxial>



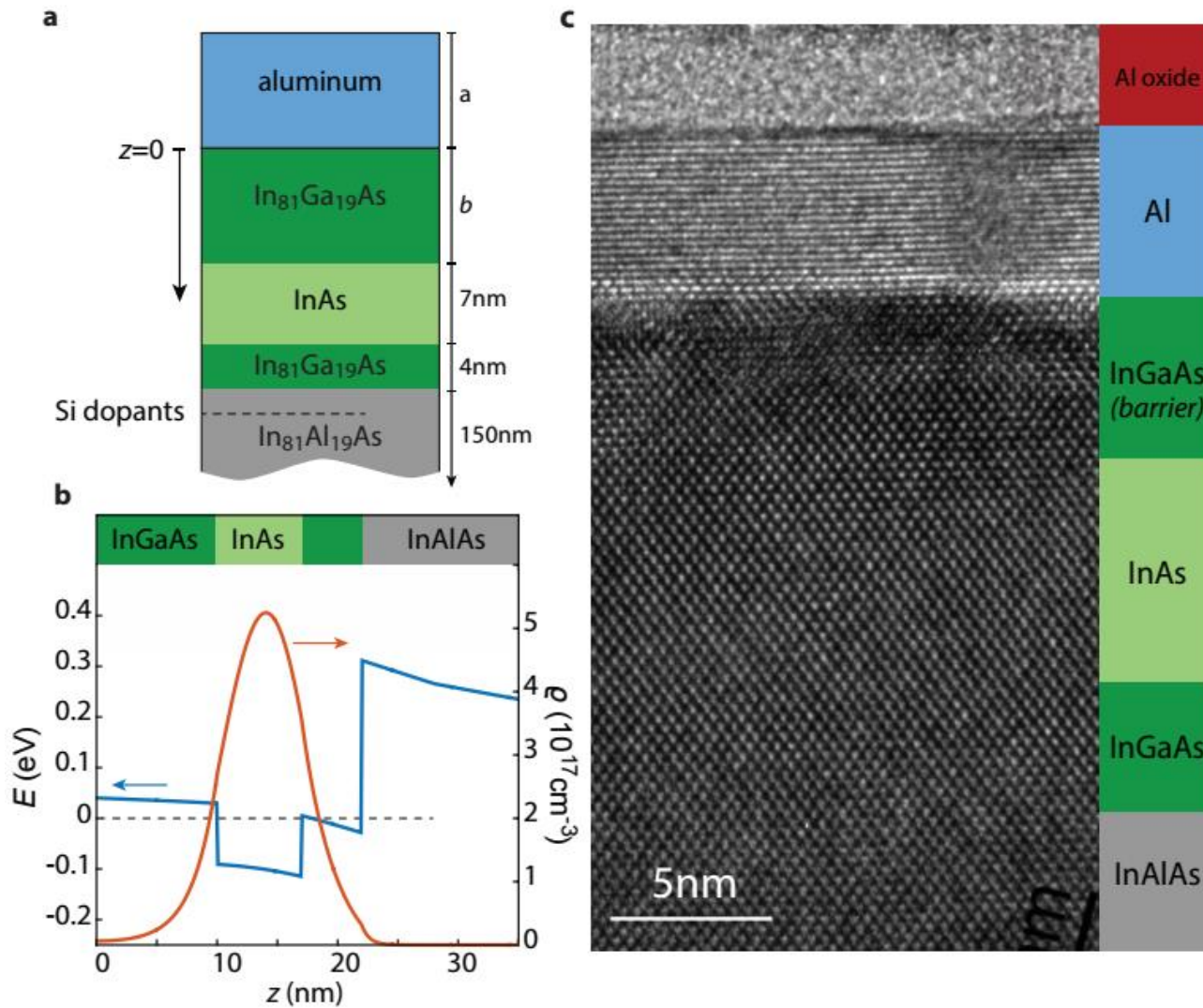
Our approach

Fabricated 1D structure from 2DEG
Epitaxial superconductor

Epitaxial Al on InAs QW



Epitaxial Al/InAs QW on InP substrate



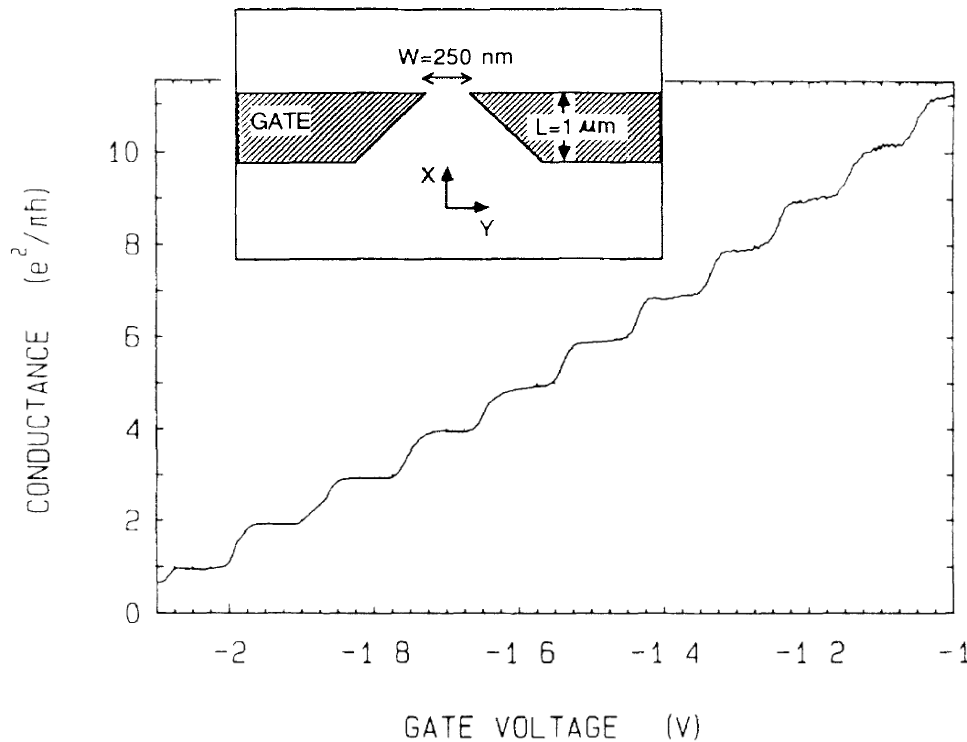
Shabani *et al.*, *PRB* **93**, 155402 (2016)

1D ballistic transport: Quantum point contact

QPC: narrow constriction in a 2DEG

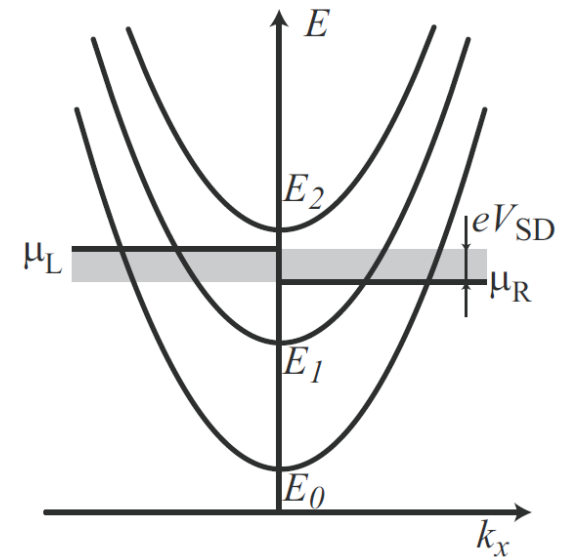
width (W) \sim Fermi wavelength (λ_F)
 $W \ll$ mean free path

Quantized conductance: $G = \frac{2e^2}{h} N$



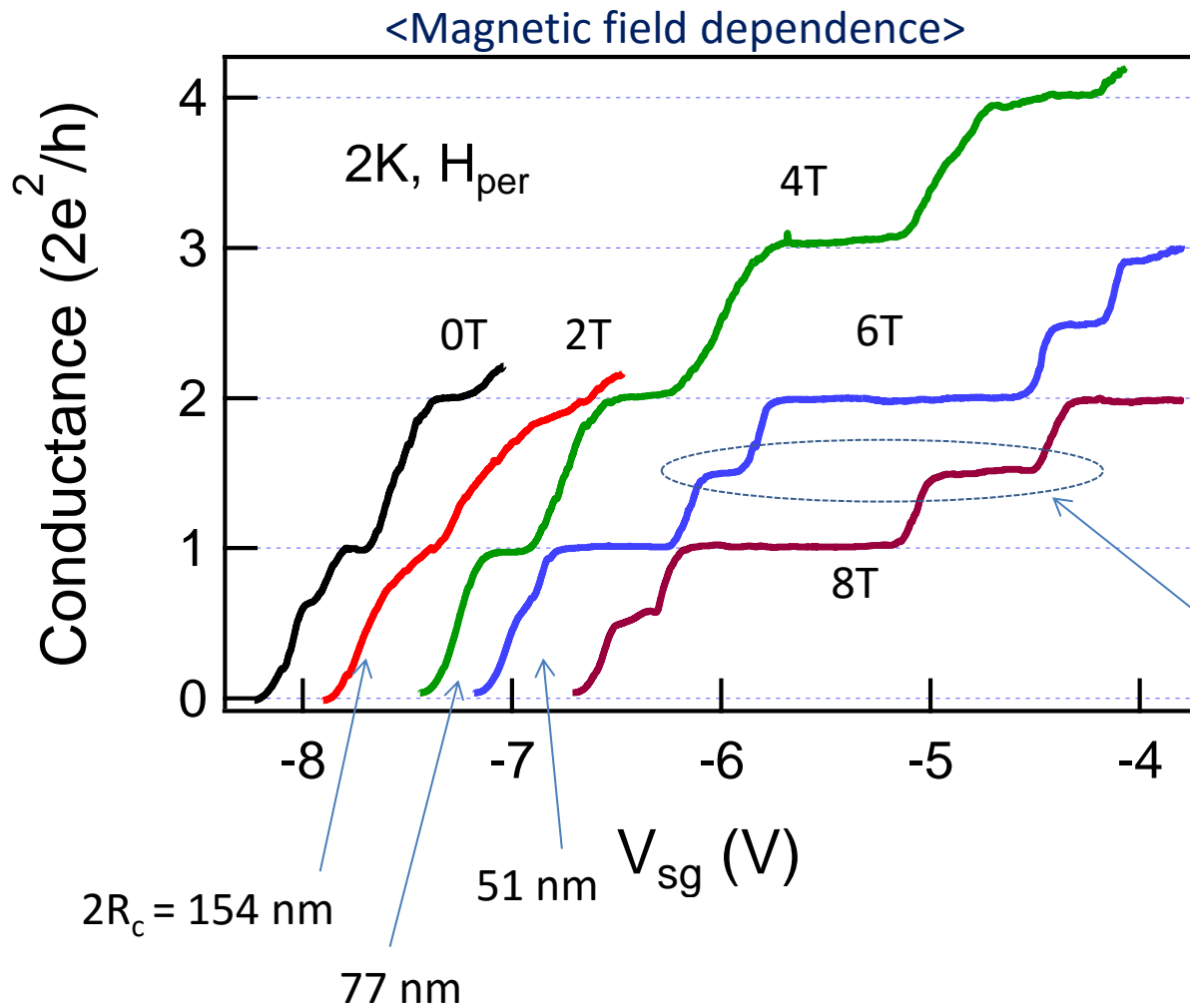
Van Wees *et al.*, *PRL* **60**, 848 (1988)

1D channel energy dispersion

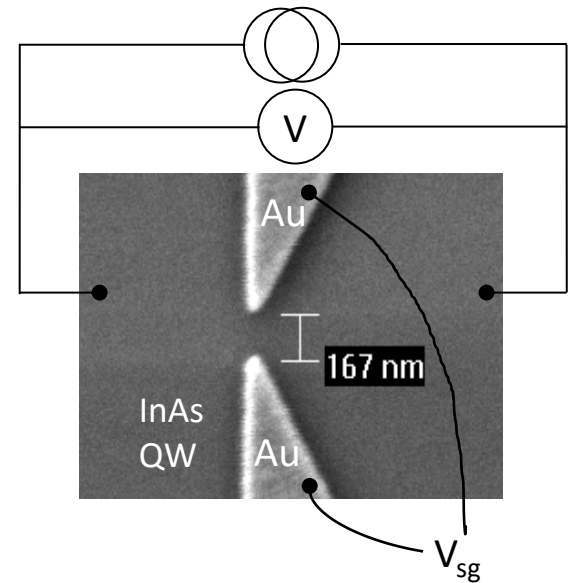


T. Ihn, *Semiconductor Nanostructures* (2010)

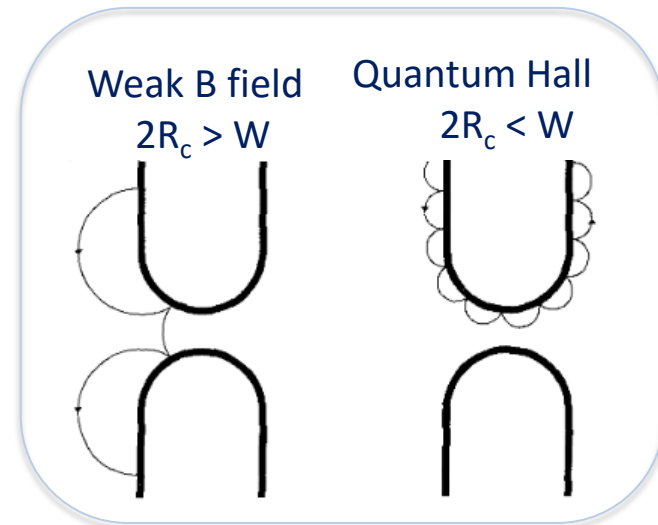
1D ballistic transport: Quantum point contact



Cyclotron radius: $R_c \equiv \hbar k_F / eB$



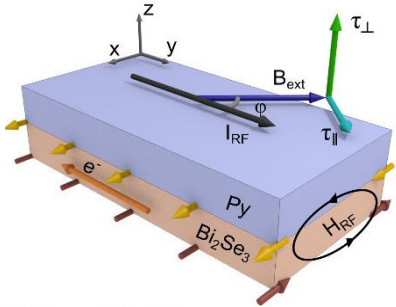
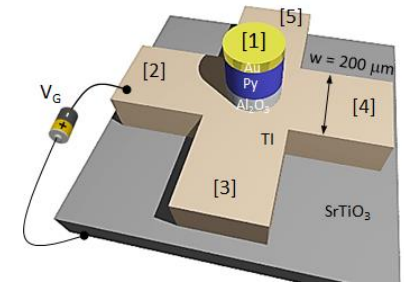
Zeeman spin splitting



Summary of the talk

1. Current-induced spin polarization of topological surface state was electrically detected.

Lee *et al.*, *Phys. Rev. B* **92**, 155312 (2015)



2. Efficient spin-charge conversion in topological insulators by spin transfer torque and spin pumping.

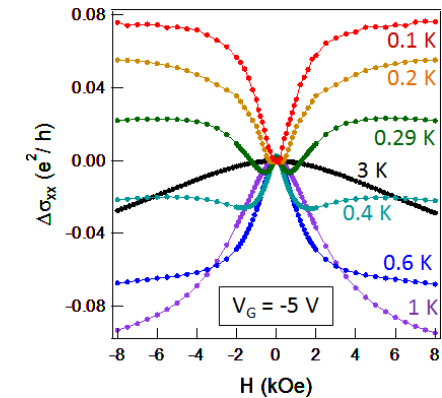
Mellnik, Lee *et al.*, *Nature* **511**, 449 (2014)

Jamali, Lee *et al.*, *Nano Lett.* **15**, 7126 (2015)

Wang, Kally, Lee *et al.*, *PRL* **117**, 076601 (2016)

3. A magnetic gap with breaking time-reversal symmetry was evidenced in topological insulator/insulating FM hybrid system.

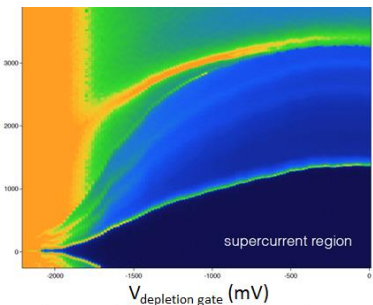
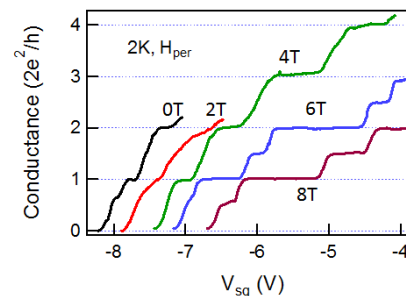
Lee *et al.* *in preparation*



4. 1D transport and superconducting proximity effect in epi-Al/InAs quantum well were demonstrate.

Lee *et al.* *in preparation*

Shojaei *et al.* *in preparation*



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Thank you for your attention!



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